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The Linear Fractional Model Theorem and Aleksandrov–Clark measures

Eva A. Gallardo-Gutiérrez and Pekka J. Nieminen

ABSTRACT

A remarkable result by Denjoy and Wolff states that every analytic self-map φ of the open unit disc $\mathbb{D}$ of the complex plane, except an elliptic automorphism, has an attractive fixed point to which the sequence of iterates $\{\varphi_n\}_{n \geq 1}$ converges uniformly on compact sets: if there is no fixed point in $\mathbb{D}$, then there is a unique boundary fixed point that does the job, called the *Denjoy–Wolff point*. This point provides a classification of the analytic self-maps of $\mathbb{D}$ into four types: maps with interior fixed point, hyperbolic maps, parabolic automorphism maps and parabolic non-automorphism maps. We determine the convergence of the Aleksandrov–Clark measures associated to maps falling in each group of such classification.

1. Introduction

Let $\mathbb{D}$ denote the open unit disc of the complex plane and φ be an analytic map taking $\mathbb{D}$ into itself. If φ has a fixed point p in $\mathbb{D}$ and is not an automorphism, then an argument based on the Schwarz Lemma shows that the sequence of iterates

$$\varphi_n = \varphi \circ \cdots \circ \varphi \quad (n \text{ times})$$

converges to p uniformly on compact subsets of $\mathbb{D}$. The Denjoy–Wolff Theorem makes a striking assertion: *If φ has no fixed point in $\mathbb{D}$, then there is still a (necessarily unique) point p in $\mathbb{T}$, the unit circle, such that $\{\varphi_n\}$ converges to p uniformly on compact subsets of $\mathbb{D}$.* This point, called the *Denjoy–Wolff point* of φ behaves like a fixed point of φ , that is, φ has radial (in fact non-tangential) limit p at p (see, for example, [16, Chapter 5]).

Moreover, it follows from the Julia–Carathéodory Theorem that the Denjoy–Wolff point is a ‘point of conformality’ in the sense of angular derivatives. More precisely, the non-tangential limit

$$\varphi'(p) = \angle \lim_{z \rightarrow p} \frac{\varphi(z) - p}{z - p},$$

exists in $(0, 1]$. Furthermore, a converse of the Denjoy–Wolff theorem also holds: *If $p \in \mathbb{T}$ is any radial-limit fixed point of a non-identity function φ and $0 < \varphi'(p) \leq 1$, then φ has no fixed points in $\mathbb{D}$ and $\{\varphi_n\}$ converges to p uniformly on compact subsets of $\mathbb{D}$, that is, p is the Denjoy–Wolff point of φ .* In such a case, the Julia–Carathéodory Theorem also asserts that

$$\varphi'(p) = \angle \lim_{z \rightarrow p} \varphi'(z) = \liminf_{z \rightarrow p} \frac{1 - |\varphi(z)|}{1 - |z|},$$

and $\varphi'(p)$ is called the *angular derivative* of φ at p .

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1.1. The Linear Fractional Model Theorem

Motivated by the restrictions on the value of the angular derivative that an analytic self-map of $\mathbb{D}$ can take at its Denjoy–Wolff point, it is possible to introduce the following general classification (see [2], Introduction).

DEFINITION 1.1. An analytic self-map φ of $\mathbb{D}$ is of

- (1) *elliptic type* if it has a fixed point in $\mathbb{D}$;
- (2) *hyperbolic type* if it has no fixed point in $\mathbb{D}$ and has angular derivative strictly less than 1 at its Denjoy–Wolff point;
- (3) *parabolic type* if it has no fixed point in $\mathbb{D}$ and has angular derivative equal to 1 at its Denjoy–Wolff point.

With this classification at hand, the *Linear Fractional Model Theorem* asserts, in some sense, that every analytic self-map φ of $\mathbb{D}$ is modelled by a linear fractional map. More precisely, let p denote the fixed point or Denjoy–Wolff point of φ . If φ is of elliptic type and $\varphi'(p) \neq 0$, then Koenigs [8] proved in 1884 that there exists a non-trivial analytic mapping σ on $\mathbb{D}$ such that Schröder’s equation

$$\sigma \circ \varphi = c\sigma$$

holds with $c = \varphi'(p)$. In 1931, Valiron [17] proved that if φ is of hyperbolic type, then the same functional equation holds for $c = 1/\varphi'(p)$ and an analytic mapping σ , taking $\mathbb{D}$ into the right half-plane and sending p to ∞ .

The parabolic case was addressed in 1979 by Pommerenke [12] and Baker and Pommerenke [1], and independently by Cowen [5], who showed that it separates into two subcases, distinguished by the behaviour of orbits relative to the pseudo-hyperbolic metric in $\mathbb{D}$, denoted here by ρ . In particular, they proved that there exists an analytic mapping σ on $\mathbb{D}$ such that either

$$\sigma \circ \varphi = \sigma + ia$$

for some real a , where the image of σ can be taken to lie in the right half-plane, or

$$\sigma \circ \varphi = \sigma + 1.$$

The former case arises when the orbits of φ are pseudo-hyperbolically separated in the sense that

$$\inf_n \rho(\varphi_{n+1}(z), \varphi_n(z)) > 0;$$

and the latter arises when they are not:

$$\inf_n \rho(\varphi_{n+1}(z), \varphi_n(z)) = 0.$$

It turns out that this separation dichotomy is independent of the base point: if it holds for one $z \in \mathbb{D}$, then it holds for all $z \in \mathbb{D}$. Thus the maps of parabolic type fall into two subclasses: *positive hyperbolic step* case (or *automorphic type*) if the orbits are separated, and *zero hyperbolic step* (or *non-automorphic type*) if they are not. The distinction between those subcases is the most subtle aspect of the Linear Fractional Model Theorem.

For a unified approach to all the cases of the Linear Fractional Model Theorem, we refer to Cowen’s paper [5] and the book [6, Section 2.4].

1.2. Aleksandrov–Clark measures

We next collect some basic facts about Aleksandrov–Clark measures. For more information on these measures, we refer the reader to [3, 9, 11, 14].

Let φ be an analytic self-map of $\mathbb{D}$. For any $\alpha \in \mathbb{T}$, an easy computation shows that the real part of the function $(\alpha + \varphi)/(\alpha - \varphi)$ is positive and harmonic in $\mathbb{D}$, so it may be expressed as the Poisson integral of a positive Borel measure τ_α supported on $\mathbb{T}$. That is,

$$\operatorname{Re} \frac{\alpha + \varphi(z)}{\alpha - \varphi(z)} = \frac{1 - |\varphi(z)|^2}{|\alpha - \varphi(z)|^2} = \int_{\mathbb{T}} P_z d\tau_\alpha,$$

where

$$P_z(\zeta) = \frac{1 - |z|^2}{|\zeta - z|^2}$$

is the Poisson kernel for $z \in \mathbb{D}$. The family of measures $\{\tau_\alpha : \alpha \in \mathbb{T}\}$ are called the *Aleksandrov–Clark measures* associated to φ . By means of the Herglotz formula, it holds that

$$\frac{\alpha + \varphi(z)}{\alpha - \varphi(z)} = \int_{\mathbb{T}} \frac{\zeta + z}{\zeta - z} d\tau_\alpha(\zeta) + ic_\alpha, \quad (1.1)$$

where

$$c_\alpha = \operatorname{Im} \frac{\alpha + \varphi(0)}{\alpha - \varphi(0)}.$$

Note that if α is given and τ is any positive and finite Borel measure on $\mathbb{T}$, then formula (1.1) can be used to construct a map φ whose Aleksandrov–Clark measure at α equals τ .

For any Borel measure τ on $\mathbb{T}$, we write $d\tau = \tau^a dm + d\tau^s$ for the Lebesgue decomposition of τ , where τ^a is the density of the absolutely continuous part, m is the normalized Lebesgue measure on $\mathbb{T}$ and τ^s is singular. It follows from the basic properties of Poisson integrals that τ_α^s is carried by the set where $\varphi(\zeta) = \alpha$ and

$$\tau_\alpha^a(\zeta) = \frac{1 - |\varphi(\zeta)|^2}{|\alpha - \varphi(\zeta)|^2}.$$

In particular, τ_α is singular if and only if φ is an inner function.

Let us recall that if the quotient $(\varphi(z) - \eta)/(z - \zeta)$ has a finite non-tangential limit at $\zeta \in \mathbb{T}$ for some $\eta \in \mathbb{T}$, then this limit is called the *angular derivative* of φ at ζ and denoted by $\varphi'(\zeta)$. It satisfies $\varphi'(\zeta) = |\varphi'(\zeta)|\bar{\zeta}\eta$ with $\eta = \varphi(\zeta)$. It turns out that the discrete part (that is, mass points, or atoms) of the Aleksandrov–Clark measures associated to φ are in correspondence with the finite angular derivatives of φ .

The map φ has a finite angular derivative at $\zeta \in \mathbb{T}$ if and only if there is $\alpha \in \mathbb{T}$ such that $\tau_\alpha(\{\zeta\}) > 0$. In that case $\varphi(\zeta) = \alpha$ and $|\varphi'(\zeta)| = \tau_\alpha(\{\zeta\})^{-1}$.

For the proof of this result, convenient references are [3, Theorem 9.2.1; 14, Theorem 3.1], where it is established in conjunction with the Julia–Carathéodory Theorem.

In view of the correspondence between the atoms of the Aleksandrov–Clark measures and angular derivatives, we can characterize the three types of analytic self-maps as follows: φ (not the identity) is of

- (1) *elliptic type* if and only if $\tau_\alpha(\{\alpha\}) < 1$ for all α ;
- (2) *hyperbolic type* if and only if $\tau_\alpha(\{\alpha\}) > 1$ for some (necessarily unique) α ;
- (3) *parabolic type* if and only if $\tau_\alpha(\{\alpha\}) = 1$ for some (necessarily unique) α .

The aim of this work is to examine the dynamical properties of the Aleksandrov–Clark measures associated to self-maps of $\mathbb{D}$ falling into this classification established by the Linear Fractional Model Theorem.

In that sense, we completely determine the convergence, in total variation norm or weak*-sense, and after a suitable normalization, of the Aleksandrov–Clark measures belonging to the iterates of φ . As it will be shown, the associated model will play a prominent role in this analysis.

Table 1 summarizes our results. By τ_α^n , we denote the Aleksandrov–Clark measure of the iterate φ_n at $\alpha \in \mathbb{T}$. In the parabolic case, we use the following terminology, which is explained in more detail in Subsection 5.2: the map φ with Denjoy–Wolff point 1 is of *finite shift* if $\sup_n \|\tau_1^n\| < \infty$ and of *infinite shift* otherwise.

2. Preliminaries

2.1. Iteration of Aleksandrov–Clark measures

In order to understand the behaviour of Aleksandrov–Clark measures under the iteration of the symbol, we will find it useful to interpret these measures from the viewpoint of composition operators. The idea is due to Sarason [15] and goes as follows: Let $\mu \in \mathcal{M}$, the space of all complex Borel measures on $\mathbb{T}$ endowed with the total variation norm. Then the Poisson integral $u(z) = \int_{\mathbb{T}} P_z d\mu$ defines a harmonic function on $\mathbb{D}$. If $\varphi: \mathbb{D} \rightarrow \mathbb{D}$ is an analytic map, then the composition $u \circ \varphi$ is also harmonic, and it is easy to see that it can be represented as the Poisson integral of a unique measure $\nu \in \mathcal{M}$. In this setting, Sarason defined $C_\varphi \mu = \nu$ and showed that the resulting linear operator C_φ acts boundedly on $\mathcal{M}$.

To describe the action of C_φ on the unit circle, note that the correspondence $C_\varphi \mu = \nu$ above can be written as

$$\begin{aligned} \int_{\mathbb{T}} P_z d\nu &= \int_{\mathbb{T}} P_{\varphi(z)} d\mu \\ &= \int_{\mathbb{T}} \left(\int_{\mathbb{T}} P_z d\tau_\alpha \right) d\mu(\alpha), \end{aligned}$$

for all $z \in \mathbb{D}$, where the second equality follows from the definition of the family $\{\tau_\alpha\}$ of Aleksandrov–Clark measures. Approximating any continuous function f on $\mathbb{T}$ by linear combinations of Poisson kernels, we arrive at the identity

$$\int_{\mathbb{T}} f d\nu = \int_{\mathbb{T}} \left(\int_{\mathbb{T}} f d\tau_\alpha \right) d\mu(\alpha).$$

TABLE 1. Convergence of Aleksandrov–Clark measures

Model for φ	Convergence of the associated Aleksandrov–Clark measures	
	φ not inner:	φ inner:
Elliptic	For any $\alpha \in \mathbb{T}$, τ_α^n is convergent in norm	For any $\alpha \in \mathbb{T}$, τ_α^n is only weak*-convergent
Hyperbolic Denjoy–Wolff Point 1 $\varphi(1) = 1$, $\varphi'(1) < 1$ $c = 1/\varphi'(1)$	For any $\alpha \in \mathbb{T} \setminus \{1\}$, τ_α^n is convergent in norm to 0 τ_1^n is NOT convergent $\frac{\tau_1^n}{c^n}$ is convergent in norm (equiv. weak*) $\Leftrightarrow \int_{\mathbb{T} \setminus \{1\}} \log \theta d\tau_1(e^{i\theta}) > -\infty$ $\frac{\tau_1^n}{\ \tau_1^n\ }$ is weak*-convergent	
Parabolic Denjoy–Wolff Point 1 $\varphi(1) = 1$, $\varphi'(1) = 1$	Finite shift case: For any $\alpha \in \mathbb{T}$, τ_α^n is convergent in norm	Infinite shift case: For any $\alpha \in \mathbb{T} \setminus \{1\}$, τ_α^n is convergent in norm to 0 τ_1^n is NOT convergent $\frac{\tau_1^n}{\ \tau_1^n\ }$ is weak*-convergent

In fact, an argument based on the monotone class theorem (cf., for example, [3, Section 9.4]) shows that this holds true for every bounded Borel function f on $\mathbb{T}$. Henceforth, we agree to write briefly

$$C_\varphi \mu = \nu = \int_{\mathbb{T}} \tau_\alpha d\mu(\alpha). \quad (2.1)$$

As a consequence of (2.1), we have the identity $\tau_\alpha = C_\varphi \delta_\alpha$ as well as the following important recursion formulas.

LEMMA 2.1. *Let φ be an analytic self-map of $\mathbb{D}$. Then, for $n \geq 2$, the Aleksandrov–Clark measure associated to φ_n at $\alpha \in \mathbb{T}$ satisfies $\tau_\alpha^n = C_\varphi \tau_\alpha^{n-1} = C_{\varphi_{n-1}} \tau_\alpha$, that is,*

$$\tau_\alpha^n = \int_{\mathbb{T}} \tau_\xi d\tau_\alpha^{n-1}(\xi) = \int_{\mathbb{T}} \tau_\xi^{n-1} d\tau_\alpha(\xi).$$

COROLLARY 2.2. *Suppose that $\tau_1(\{1\}) = c > 0$, or equivalently $\varphi(1) = 1$ and $\varphi'(1) = 1/c$. Then, for $n \geq 2$,*

$$\tau_1^n = c\tau_1^{n-1} + \nu_n \quad \text{with} \quad \nu_n = \int_{\mathbb{T} \setminus \{1\}} \tau_\xi^{n-1} d\tau_1(\xi).$$

Hence, τ_1^n can be expressed as the following sum of positive measures:

$$\tau_1^n = c^n \delta_1 + c^{n-1} \nu_1 + c^{n-2} \nu_2 + \cdots + \nu_n,$$

where ν_1 is the restriction of τ_1 to $\mathbb{T} \setminus \{1\}$.

REMARK 2.3. It is of interest to note that the singular parts of τ_1^{n-1} and ν_n above have disjoint supports in the sense of measure. Indeed, if

$$E_{n-1} = \{\zeta \in \mathbb{T} : \varphi_{n-1}(\zeta) = 1\},$$

then $(\tau_1^{n-1})^s(\mathbb{T} \setminus E_{n-1}) = 0$, while $(\tau_\xi^{n-1})^s(E_{n-1}) = 0$ for each $\xi \neq 1$ and hence $\nu_n^s(E_{n-1}) = 0$. As a consequence, the measures arising in the above expansion for τ_1^n have singular parts with disjoint supports.

2.2. Half-plane model for hyperbolic and parabolic types

Assume that φ is of either hyperbolic or parabolic type with Denjoy–Wolff point 1 and $\varphi'(1) = 1/c \in (0, 1]$. This means that $\tau_1(\{1\}) = c \geq 1$ and hence

$$\frac{1 + \varphi(z)}{1 - \varphi(z)} = c \frac{1 + z}{1 - z} + ia + \int_{\mathbb{T} \setminus \{1\}} \frac{\zeta + z}{\zeta - z} d\tau_1(\zeta),$$

where $a = \text{Im}(1 + \varphi(0))/(1 - \varphi(0))$.

We often transform this into the right half-plane $\mathbb{H}$ by sending the Denjoy–Wolff point 1 to ∞ via the Cayley transformations

$$w = \frac{1 + z}{1 - z}, \quad z = \frac{w - 1}{w + 1}, \quad \Phi(w) = \frac{1 + \varphi(z)}{1 - \varphi(z)}, \quad \varphi(z) = \frac{\Phi(w) - 1}{\Phi(w) + 1}. \quad (2.2)$$

Then Φ maps $\mathbb{H}$ into itself and is of the form

$$\Phi(w) = cw + ia + G(w) = cw + \Gamma(w),$$

where G and $\Gamma = ia + G$ have the property that

$$\angle \lim_{w \rightarrow \infty} \frac{G(w)}{w} = \angle \lim_{w \rightarrow \infty} \frac{\Gamma(w)}{w} = 0.$$

Here w approaches ∞ non-tangentially, that is, within a Stolz angle $\{w : |\text{Im } w| < \beta \text{Re } w\}$ for some $\beta > 0$. Thus Φ has ∞ as its Denjoy–Wolff point with the angular derivative $\Phi'(\infty) = c$. (See [2, Section 4].)

If necessary, the integral defining G above can be transformed into an integral calculated over the real line. More precisely, writing $w = x + iy$ with $x > 0$ we have

$$G(w) = \int_{\mathbb{R}} \frac{i\eta w - 1}{i\eta - w} d\tilde{\tau}_1(\eta) = U(w) + iV(w),$$

where

$$U(w) = \int_{\mathbb{R}} \frac{x}{x^2 + (\eta - y)^2} (1 + \eta^2) d\tilde{\tau}_1(\eta),$$

$$V(w) = \int_{\mathbb{R}} \left\{ -\eta + \frac{(\eta - y)(1 + \eta^2)}{x^2 + (\eta - y)^2} \right\} d\tilde{\tau}_1(\eta),$$

and $\tilde{\tau}_1$ is the pull-back of τ_1 to $\mathbb{R}$, that is,

$$d\tilde{\tau}_1(\eta) = d\tau_1 \left(\frac{i\eta - 1}{i\eta + 1} \right).$$

It should be noted that here $\tilde{\tau}_1$ can be any positive Borel measure with finite mass on the real line. The integrals defining U and V make sense because the integrands are bounded functions of η for each fixed w . However, the integral $\int \eta d\tilde{\tau}_1(\eta)$ is not necessarily well defined, and hence a different normalization for V cannot always be used (cf. the discussion in [7, p. 105]).

3. Warm-up: the elliptic case

If φ is not an automorphism and fixes a point $p \in \mathbb{D}$ (that is, φ is of elliptic type), then φ_n converges to p uniformly on compact sets and therefore for each $\alpha \in \mathbb{T}$,

$$\lim_{n \rightarrow \infty} \int_{\mathbb{T}} P_z d\tau_{\alpha}^n = \lim_{n \rightarrow \infty} \frac{1 - |\varphi_n(z)|^2}{|\alpha - \varphi_n(z)|^2} = \frac{1 - |p|^2}{|\alpha - p|^2}$$

for all $z \in \mathbb{D}$. Hence, the sequence $\{\tau_{\alpha}^n\}$ converges in the weak*-topology to the measure $((1 - |p|^2)/|\alpha - p|^2)m$. One may ask if $\{\tau_{\alpha}^n\}$ converges in the norm of $\mathcal{M}$. Having in mind that the Aleksandrov–Clark measures associated to inner functions are singular, one easily deduces the next result which says that it is not always the case.

PROPOSITION 3.1. *Let φ be a non-automorphic inner function fixing $p \in \mathbb{D}$. For any $\alpha \in \mathbb{T}$, the sequence of Aleksandrov–Clark measures $\{\tau_{\alpha}^n\}$ associated to φ_n at α does not converge in norm.*

Nevertheless, if φ is not inner, then the answer is positive.

PROPOSITION 3.2. *Let φ be an analytic self-map of $\mathbb{D}$ fixing $p \in \mathbb{D}$. Assume that φ is not an inner function. Then, for any $\alpha \in \mathbb{T}$, the sequence of Aleksandrov–Clark measures $\{\tau_{\alpha}^n\}$ associated to φ_n at α converges in norm to $((1 - |p|^2)/|\alpha - p|^2)m$.*

Proof. We first assume that $\varphi(0) = 0$ and show that $\{\tau_{\alpha}^n\}$ converges to m in norm. Let $\mathcal{M}_0$ denote the subspace of $\mathcal{M}$ consisting of measures μ such that $\mu(\mathbb{T}) = 0$; or equivalently $P[\mu](0) = 0$ where $P[\mu]$ is the Poisson integral of μ :

$$P[\mu](z) = \int_{\mathbb{T}} P_z(\zeta) d\mu(\zeta).$$

Since $\varphi(0) = 0$, C_{φ} is taking $\mathcal{M}_0$ into itself. Moreover, having in mind that $C_{\varphi_n} \delta_{\alpha} = \tau_{\alpha}^n$ for every $n \geq 1$ as well as trivially $C_{\varphi} m = m$, one deduces that

$$\|\tau_{\alpha}^n - m\| = \|C_{\varphi_n}(\delta_{\alpha} - m)\| \leq \|C_{\varphi}| \mathcal{M}_0\|^n \|\delta_{\alpha} - m\|.$$

It is enough to check that $\|C_\varphi|\mathcal{M}_0\| < 1$. Since φ is non-inner, there exists a set $E \subset \mathbb{T}$ of positive Lebesgue measure such that, for each $\zeta \in E$, one has $|\varphi(\zeta)| < 1$ and hence

$$g(\zeta) = \inf_{\xi \in \mathbb{T}} \tau_\xi^a(\zeta) = \inf_{\xi \in \mathbb{T}} \frac{1 - |\varphi(\zeta)|^2}{|\xi - \varphi(\zeta)|^2} > 0.$$

Put $g = 0$ on $\mathbb{T} \setminus E$. Then, for $\mu \in \mathcal{M}_0$, we have (in the sense of (2.1))

$$C_\varphi \mu = \int_{\mathbb{T}} \tau_\xi d\mu(\xi) = \int_{\mathbb{T}} (\tau_\xi - g \cdot m) d\mu(\xi).$$

Note that each of the measures $\tau_\xi - g \cdot m$ is positive. Hence,

$$\|C_\varphi|\mathcal{M}_0\| \leq \sup_{\xi \in \mathbb{T}} \|\tau_\xi - g \cdot m\| = 1 - \int_{\mathbb{T}} g dm < 1$$

because $\|\tau_\xi\| = 1$. This completes the proof in the case when $\varphi(0) = 0$.

Now assume that $p \neq 0$ and define $\psi = \sigma_p \circ \varphi \circ \sigma_p$ where σ_p is the disc automorphism that exchanges p and 0. Then $\psi(0) = 0$ and $\varphi_n = \sigma_p \circ \psi_n \circ \sigma_p$ for each n . Hence, $\tau_\alpha^n = C_{\varphi_n} \delta_\alpha = C_{\sigma_p} C_{\psi_n} C_{\sigma_p} \delta_\alpha$. Here it is easy to verify that

$$C_{\sigma_p} \delta_\alpha = \frac{1 - |p|^2}{|\alpha - p|^2} \delta_{\sigma_p(\alpha)}.$$

This completes the proof since $C_{\sigma_p} C_{\psi_n} \delta_{\sigma_p(\alpha)} \rightarrow C_{\sigma_p} m = m$ in norm by the first part of the proof. $\square$

4. The hyperbolic case

In this section, we will assume that φ is an analytic self-map of $\mathbb{D}$ of hyperbolic type, that is, φ has no fixed point in $\mathbb{D}$ and the angular derivative at its Denjoy–Wolff point is strictly less than 1. Without loss of generality, we will assume throughout this section that 1 is the Denjoy–Wolff point of φ .

Our first statement concerns the sequence of Aleksandrov–Clark measures associated to a point $\alpha \in \mathbb{T} \setminus \{1\}$. It is based on the fact that the norm of τ_α^n is given by

$$\|\tau_\alpha^n\| = \frac{1 - |\varphi_n(0)|^2}{|\alpha - \varphi_n(0)|^2}.$$

PROPOSITION 4.1. *Let φ be an analytic self-map of $\mathbb{D}$ with Denjoy–Wolff point 1. For any $\alpha \in \mathbb{T} \setminus \{1\}$, the sequence of Aleksandrov–Clark measures $\{\tau_\alpha^n\}$ converges to 0 in norm.*

Therefore, in both parabolic and hyperbolic types, we are reduced to study the convergence of the sequence of measures $\{\tau_1^n\}$.

4.1. The model

If φ is of hyperbolic type with Denjoy–Wolff point 1 and $c = 1/\varphi'(1)$, then $\tau_1(\{1\}) = c > 1$. The half-plane model for φ is of the form

$$\Phi(w) = cw + \Gamma(w),$$

where $\Gamma: \mathbb{H} \rightarrow \mathbb{H}$ is analytic with angular derivative 0 at infinity, that is, $\Gamma(w)/w \rightarrow 0$ as $w \rightarrow \infty$ non-tangentially (see Subsection 2.2). Moreover, by Corollary 2.2, the Aleksandrov–Clark measures of the iterates φ_n at 1 have the expansion

$$\tau_1^n = c^n \delta_1 + c^{n-1} \nu_1 + c^{n-2} \nu_2 + \cdots + \nu_n.$$

Obviously, $\|\tau_1^n\| \geq c^n$ and hence $\|\tau_1^n\| \rightarrow \infty$ as $n \rightarrow \infty$. So, in order to discuss any sort of convergence, we have to normalize the sequence $\{\tau_1^n\}$. As we will see, the results depend strongly on the normalization method.

4.2. A function-theoretic normalization

We first study the normalized sequence of measures

$$\varphi'_n(1)\tau_1^n = \frac{\tau_1^n}{c^n}$$

for $n \geq 1$. This normalization goes back to classical works of Koenigs, Wolff and Valiron, and from a function-theoretic perspective can be considered ‘a natural one’ since it corresponds to the maps $\varphi_n/\varphi'_n(1)$. However, it does not always work: it is not always the case that

$$\sup_n \frac{\|\tau_1^n\|}{c^n} = \sup_n \frac{\operatorname{Re} \Phi_n(1)}{c^n} < \infty; \quad (4.1)$$

for instance, see the discussion by Valiron [17] (starting on p. 120), where the mapping

$$\Phi(w) = 2w + \frac{w}{\log(3+w)} \quad (4.2)$$

provides a counter-example to (4.1).

In fact, it was later shown by Pommerenke [13] that (4.1) holds if and only if

$$\int_1^\infty \frac{|\Phi(t) - ct|}{t^2} dt = \int_1^\infty \frac{|\Gamma(t)|}{t^2} dt < \infty. \quad (4.3)$$

This can be viewed as a regularity requirement imposed on φ at its Denjoy–Wolff point; for instance, it is seen to be fulfilled if φ satisfies an angular smoothness condition of order $1 + \varepsilon$ at 1 for some $\varepsilon > 0$, which in terms of the half-plane model means that $\Gamma(w)/|w|^{1-\varepsilon} \rightarrow 0$ as $w \rightarrow \infty$ non-tangentially (cf., for example, [2, p. 51]).

Next result establishes that the condition (4.1) is equivalent to a sort of regularity of the measure τ_1 .

PROPOSITION 4.2. *Suppose that the Aleksandrov–Clark measure τ_1 of φ satisfies $\tau_1(\{1\}) = c > 1$. Then $\sup_n \|\tau_1^n\|/c^n < \infty$ if and only if*

$$\int_{\mathbb{T} \setminus \{1\}} \log |\theta| d\tau_1(e^{i\theta}) > -\infty.$$

Proof. We prove that the proposed condition is equivalent to Pommerenke’s condition (4.3). After performing the change of variables $t = (1+r)/(1-r)$, we see that (4.3) is equivalent to the pair of conditions

$$\int_0^1 \operatorname{Re} \gamma(r) dr < \infty \quad \text{and} \quad \int_0^1 |\operatorname{Im} \gamma(r)| dr < \infty, \quad (4.4)$$

where (see Subsection 2.2)

$$\gamma(r) = \int_{\mathbb{T} \setminus \{1\}} \frac{e^{i\theta} + r}{e^{i\theta} - r} d\tau_1(e^{i\theta}) = \int_{\mathbb{T} \setminus \{1\}} P_r(e^{i\theta}) d\tau_1(e^{i\theta}) + i \int_{\mathbb{T} \setminus \{1\}} \frac{2r \sin \theta}{|e^{i\theta} - r|^2} d\tau_1(e^{i\theta}).$$

Note that in the denominator we may estimate $|e^{i\theta} - r|^2 \simeq (1-r)^2 + \theta^2$ for $|\theta| \leq \pi$ (meaning that the two functions are bounded by a constant multiple of each other). Indeed, one may write $|e^{i\theta} - r|^2 = (1-r)^2 + 2r(1 - \cos \theta)$ and proceed by using the fact that $1 - \cos \theta \simeq \theta^2$.

Consequently, for the Poisson kernel we have the estimate

$$\int_0^1 P_r(e^{i\theta}) dr \simeq \int_0^1 \frac{(1-r) dr}{(1-r)^2 + \theta^2} = \frac{1}{2} \log(1 + \theta^2) - \log |\theta|,$$

whenever $|\theta| \leq \pi$. This, along with an application of Fubini's theorem, shows that the first part of (4.4) is equivalent to the condition of the proposition; observe that the function $\frac{1}{2} \log(1 + \theta^2)$ is bounded and hence plays no role in integrability. On the other hand, since

$$\int_0^1 \frac{|2r \sin \theta|}{|e^{i\theta} - r|^2} dr \leq C \int_0^1 \frac{|\theta| dr}{(1-r)^2 + \theta^2} \leq C \int_0^\infty \frac{ds}{s^2 + 1} < \infty$$

with $C > 0$ being a constant independent of θ for $|\theta| \leq \pi$, we see that the second part of (4.4) is always satisfied. This completes the proof. $\square$

THEOREM 4.3. *Let φ be of hyperbolic type with Denjoy–Wolff point 1 and $\varphi'(1) = 1/c$. If $\sup_n \|\tau_1^n\|/c^n < \infty$, then the sequence $\{\tau_1^n/c^n\}$ converges in norm to a positive measure $\mu \in \mathcal{M}$, satisfying*

$$c\mu = \int_{\mathbb{T}} \tau_\xi d\mu(\xi) \quad (4.5)$$

and $\mu(\{1\}) = 1$.

Proof. Since the measures $\tau_1^{n+1}/c^{n+1} - \tau_1^n/c^n$ are positive for all n , the existence of the limit measure follows from the boundedness assumption by the completeness of $\mathcal{M}$. By Lemma 2.1, we have the identity

$$c \frac{\tau_1^{n+1}}{c^{n+1}} = \int_{\mathbb{T}} \tau_\xi \frac{d\tau_1^n(\xi)}{c^n},$$

and so formula (4.5) is obtained by passing to the limit as $n \rightarrow \infty$. $\square$

Relation (4.5) gives rise to Schröder's functional equation $c\sigma = \sigma \circ \varphi$ in a straightforward way. Indeed, in terms of composition operators (see Subsection 2.1), (4.5) says that $c\mu = C_\varphi\mu$. So, if $\sigma: \mathbb{D} \rightarrow \mathbb{H}$ is defined as the Herglotz integral of μ , that is,

$$\sigma(z) = \int_{\mathbb{T}} \frac{\zeta + z}{\zeta - z} d\mu(\zeta), \quad (4.6)$$

then we must have $c\operatorname{Re} \sigma = \operatorname{Re} \sigma \circ \varphi$. Schröder's equation is then obtained by adding a suitable imaginary constant to σ .

Next example provides a map φ of hyperbolic type whose Aleksandrov–Clark measure fails the logarithmic integrability condition of Proposition 4.2. Such a map can be seen as an alternative to Valiron's example (4.2), showing that 'the natural normalization' Φ_n/c^n does not always yield a finite limit as $n \rightarrow \infty$.

EXAMPLE 4.4. We define φ in the form

$$\frac{1 + \varphi(z)}{1 - \varphi(z)} = 2 \frac{1 + z}{1 - z} + \int_{\mathbb{T}} \frac{\zeta + z}{\zeta - z} d\nu(\zeta),$$

where ν is the discrete measure $\nu = \sum_{k=1}^\infty k^{-2} \delta_{\zeta_k}$ with $\zeta_k = \exp(2^{-k}i)$. Then the Aleksandrov–Clark measure associated to φ at 1 is $\tau_1 = 2\delta_1 + \nu$. In particular, φ is of hyperbolic type with Denjoy–Wolff point 1 and $\varphi'(1) = \frac{1}{2}$. A simple computation shows that

$$\int_{\mathbb{T}} \log |\theta| d\nu(e^{i\theta}) = \sum_k (\log 2^{-k}) k^{-2} = -\infty.$$

4.3. A stronger normalization

Let again φ be of hyperbolic type with Denjoy–Wolff point 1, and consider the probability measures

$$\mu_n = \frac{\tau_1^n}{\|\tau_1^n\|} = \frac{|1 - \varphi_n(0)|^2}{1 - |\varphi_n(0)|^2} \tau_1^n.$$

From a purely measure-theoretic perspective, this might appear to be a more appropriate normalization than the one considered above.

It can be shown that the sequence $\{\mu_n\}$ always converges in weak*-sense to a non-trivial measure μ . Indeed, in terms of the Cayley transforms (2.2) we may write

$$\int_{\mathbb{T}} P_z d\mu_n = \frac{\operatorname{Re} \Phi_n(w)}{\operatorname{Re} \Phi_n(1)},$$

and it is known that this quotient converges to a finite limit for each $w \in \mathbb{H}$ (see, for example, [12]). Moreover, the intertwining relation (4.5) still holds. This is seen by letting $n \rightarrow \infty$ in the identity

$$\frac{\operatorname{Re} \Phi_{n+1}(1)}{\operatorname{Re} \Phi_n(1)} \mu_{n+1} = \int_{\mathbb{T}} \tau_\xi d\mu_n(\xi)$$

and noting that the quotient $\operatorname{Re} \Phi_{n+1}(1)/\operatorname{Re} \Phi_n(1) = \operatorname{Re} \Phi(\Phi_n(1))/\operatorname{Re} \Phi_n(1)$ tends to c by the Julia–Carathéodory Theorem since $\Phi_n(1)$ is known to converge to ∞ non-tangentially (see, for example, [16, Proposition 5.2]).

5. The parabolic case

In this section, we will assume that φ is an analytic self-map of $\mathbb{D}$ of parabolic type, that is, φ has no fixed point in $\mathbb{D}$ and the angular derivative at its Denjoy–Wolff point is 1. Without loss of generality, we will assume throughout this section that 1 is the Denjoy–Wolff point of φ .

As we already pointed out in Proposition 4.1, for any $\alpha \in \mathbb{T} \setminus \{1\}$, the sequence of Aleksandrov–Clark measures $\{\tau_\alpha^n\}$ converges to zero in the norm of $\mathcal{M}$. Therefore, we will focus on the convergence of $\{\tau_1^n\}$.

5.1. The model

If φ is of parabolic type with Denjoy–Wolff point 1, then $\tau_1(\{1\}) = 1$. As in the hyperbolic case (see Subsection 4.1), the half-plane model of φ is of the form

$$\Phi(w) = w + \Gamma(w),$$

where $\Gamma(w)/w \rightarrow 0$ as $w \rightarrow \infty$ non-tangentially, and the Aleksandrov–Clark measure of φ_n at 1 has the expression given by Corollary 2.2:

$$\tau_1^n = \delta_1 + \nu_1 + \nu_2 + \cdots + \nu_n.$$

5.2. Finite vs. infinite shift

From a measure-theoretic point of view, two natural cases now arise in a way which is similar to the hyperbolic setting in the previous section. Recalling that

$$\|\tau_1^n\| = \operatorname{Re} \Phi_n(1) = \frac{1 - |\varphi_n(0)|^2}{|1 - \varphi_n(0)|^2}$$

we may say (as in [4]) that the map φ (or Φ) has

(1) *finite shift* if

$$\sup_n \|\tau_1^n\| = \sup_n \operatorname{Re} \Phi_n(1) < \infty;$$

(2) *infinite shift* otherwise.

We remark that the value $\Phi_n(1)$ here could be replaced by $\Phi_n(w)$ for any $w \in \mathbb{H}$ (see [4, Proposition 3.2]).

Next result follows similarly to Theorem 4.3, just taking $c = 1$.

THEOREM 5.1. *Let φ be of parabolic type with Denjoy–Wolff point 1. If φ has finite shift, then the sequence $\{\tau_1^n\}$ converges in norm to a positive measure $\mu \in \mathcal{M}$, satisfying*

$$\mu = \int_{\mathbb{T}} \tau_\xi d\mu(\xi) \quad (5.1)$$

and $\mu(\{1\}) = 1$.

We observe that if $\sigma: \mathbb{D} \rightarrow \mathbb{H}$ is the Herglotz integral of μ above (cf. equation (4.6)), then relation (5.1) yields the functional equation $\sigma + ia = \sigma \circ \varphi$. Moreover, the convergence $\tau_1^n \rightarrow \mu$ implies that the associated Herglotz integrals converge uniformly on compact sets:

$$\frac{1 + \varphi_n(z)}{1 - \varphi_n(z)} - y_n = \Phi_n \left(\frac{1+z}{1-z} \right) - y_n \longrightarrow \sigma(z), \quad (5.2)$$

where $y_n = \operatorname{Im}(1 + \varphi_n(0))/(1 - \varphi_n(0)) = \operatorname{Im} \Phi_n(1)$. It is worth noting that the intertwining map σ thus obtained is non-constant; in fact, since $\mu(\{1\}) = 1$, the map $(\sigma - 1)/(\sigma + 1)$, which takes $\mathbb{D}$ into itself, has angular derivative 1 at 1.

The case of infinite shift is much more delicate. Recall from Section 1 that the parabolic self-maps of the disc in general fall into two fundamental subtypes, those with positive or zero hyperbolic step, according to whether the orbits of the map are separated or not relative to the pseudo-hyperbolic metric. Thus, combining the concepts of positive/zero hyperbolic type and finite/infinite shift, there is in principle a classification of parabolic maps into four different classes. Nevertheless, as observed by Poggi-Corradini [10, Proposition 4.1] (see also [4, Proposition 3.3]), maps with finite shift always have positive hyperbolic step.

Instead of the normalization in (5.2), Pommerenke [12] used the normalization

$$\sigma_n(z) = \frac{1}{x_n} \left[\Phi_n \left(\frac{1+z}{1-z} \right) - y_n \right],$$

where $x_n + iy_n = \Phi_n(1)$. He showed that then the sequence $\{\sigma_n\}$ always converges to an analytic map $\sigma: \mathbb{D} \rightarrow \mathbb{H}$ uniformly on compact sets. In addition:

- (1) in the positive hyperbolic step case $\sigma \circ \varphi = \sigma + ia$ for some real $a \neq 0$ (whence σ is non-constant);
- (2) in the zero hyperbolic step case $\sigma \equiv 1$.

Since $x_n = \|\tau_1^n\|$, we deduce the following result.

PROPOSITION 5.2. *Let φ be of parabolic type with Denjoy–Wolff point 1. Then the sequence of measures $\{\tau_1^n/\|\tau_1^n\|\}$ is weak*-convergent to a positive measure $\mu \in \mathcal{M}$ which satisfies (5.1).*

REMARK 5.3. In the case of zero hyperbolic step, the measure μ above is just the Lebesgue measure m because $\sigma \equiv 1$. To obtain a non-trivial intertwining map, a somewhat trickier normalization can be used, as discovered by Baker and Pommerenke [1]. However, it seems not so straightforward to interpret it from the viewpoint of Aleksandrov–Clark measures.

We also point out that the properties of positive/zero hyperbolic type and finite/infinite shift are not directly determined by regularity properties of the map φ (as in the hyperbolic case). However, under certain smoothness assumptions placed on φ at its Denjoy–Wolff point, finite shift becomes equivalent to positive hyperbolic step and distinguishing between the two remaining classes becomes easier. See [2, 4] for results in this direction.

We end this section with a couple of examples.

EXAMPLE 5.4. Let $\tau_1 = \delta_1 + a\delta_\alpha$ with $\alpha \in \mathbb{T} \setminus \{1\}$ and $a > 0$, and let $b \in \mathbb{R}$. Define $\varphi: \mathbb{D} \rightarrow \mathbb{D}$ by the formula

$$\frac{1 + \varphi(z)}{1 - \varphi(z)} = \frac{1 + z}{1 - z} + a \frac{\alpha + z}{\alpha - z} + ib,$$

so that τ_1 is the Aleksandrov–Clark measure of φ at 1. Equivalently, in the right half-plane

$$\Phi(w) = w + a \frac{i\eta w - 1}{i\eta - w} + ib = w + a \frac{1 + \eta^2}{w - i\eta} + i(b - a\eta),$$

where $i\eta = (1 + \alpha)/(1 - \alpha)$ is a point of the imaginary axis. Therefore, $\varphi''(1) = i(b - a\eta)$, and [2, Theorem 4.4] shows that φ has positive hyperbolic step if and only if $b \neq a\eta$. In this case, φ is also of finite shift. Note that the hyperbolic step and shift of φ depend on all the parameters a , α (or η) and b . In the symmetric case $\alpha = -1$ (or $\eta = 0$), however, the value of a has no effect.

REMARK 5.5. In the case of zero hyperbolic step, the preceding example shows that the convergence given by Proposition 5.2 cannot be in norm in general. Indeed, all the measures τ_1^n there are discrete, though the limiting measure is the Lebesgue measure, as observed in Remark 5.3.

PROPOSITION 5.6. Let φ be of parabolic type with Denjoy–Wolff point 1; that is, $\tau_1(\{1\}) = 1$. Suppose that the absolutely continuous part of τ_1 satisfies $d\tau_1^a = g dm$, where $0 \leq g \in L^1(\mathbb{T})$ and g is continuous at 1 with $g(1) > 0$. Then $\|\tau_1^n\| \rightarrow \infty$ as $n \rightarrow \infty$, so φ is of infinite shift.

Proof. Recall that $\tau_1^n = \delta_1 + \nu_1 + \nu_2 + \cdots + \nu_n$ with

$$\nu_n = \int_{\mathbb{T} \setminus \{1\}} \tau_\xi^{n-1} d\tau_1(\xi).$$

Now we have

$$\nu_n(\mathbb{T}) = \int_{\mathbb{T} \setminus \{1\}} \tau_\xi^{n-1}(\mathbb{T}) d\tau_1(\xi) \geq \int_{\mathbb{T}} \frac{1 - |\varphi_{n-1}(0)|^2}{|\xi - \varphi_{n-1}(0)|^2} g(\xi) dm(\xi) \longrightarrow g(1) > 0$$

since $\varphi_{n-1}(0) \rightarrow 1$. Therefore, $\|\tau_1^n\| = \tau_1^n(\mathbb{T}) \rightarrow \infty$. $\square$

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