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Observation of New Charmless Decays of Bottom Hadrons

We search for new charmless decays of neutral $b$–hadrons to pairs of charged hadrons with the upgraded Collider Detector at the Fermilab Tevatron. Using a data sample corresponding to $1 \text{fb}^{-1}$ of integrated luminosity, we report the first observation of the $B^0 \rightarrow K^- \pi^+$ decay, with a significance of $8.2 \sigma$, and measure $B(B^0 \rightarrow K^- \pi^+) = (5.0 \pm 0.7 \text{ (stat)} \pm 0.5 \text{ (syst)}) \times 10^{-6}$. We also report the first observation of charmless $b$–baryon decays in the channels $\Lambda_b^0 \rightarrow p\pi^-$ and $\Lambda_b^0 \rightarrow pK^-$ with significances of $6.0 \sigma$ and $11.5 \sigma$ respectively, and we measure $B(\Lambda_b^0 \rightarrow p\pi^-) = (3.5 \pm 0.6 \text{ (stat)} \pm 0.9 \text{ (syst)}) \times 10^{-6}$ and $B(\Lambda_b^0 \rightarrow pK^-) = (5.6 \pm 0.8 \text{ (stat)} \pm 1.5 \text{ (syst)}) \times 10^{-6}$. No evidence is found for the decays $B^0 \rightarrow K^+K^-$ and $B^0 \rightarrow \pi^+\pi^-$, and we set an improved upper limit $B(B^0 \rightarrow \pi^+\pi^-) < 1.2 \times 10^{-5}$ at the 90% confidence level. All quoted branching fractions are measured using $B(B^0 \rightarrow K^-\pi^+)$ as a reference.

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Two-body non-leptonic charmless decays of $b$–hadrons are among the most widely studied processes in flavor physics. The variety of open channels involving similar final states provides crucial experimental information to improve the accuracy of effective models of strong interaction dynamics. The quark-level transition $b \rightarrow u$ makes decay amplitudes sensitive to $\gamma$, the least known angle of the quark-mixing (Cabibbo-Kobayashi-Maskawa, CKM) matrix. Significant contributions from higher-order (‘penguin’) transitions provide sensitivity to the possible presence of new physics in internal loops, if the CKM angle $\gamma$ is of particular interest, because its branching fraction is sensitive to the CKM angle $\gamma$. The current experimental bound on $\gamma$ is lower than most predictions.

Of the possible $B_s^0$ decay modes into pairs of charmless pseudoscalar mesons, only the $B_s^0 \rightarrow K^+ K^-$ has been observed to date. The $B_s^0 \rightarrow K^- \pi^+$ is of particular interest, because its branching fraction is sensitive to the CKM angle $\gamma$ and the current experimental bound on $\gamma$ is lower than most predictions.

A measurement of the branching fraction of the $B_s^0 \rightarrow \pi^+ \pi^-$ mode, along with the $B_s^0 \rightarrow K^+ K^-$ mode, would allow a determination of the strength of penguin-annihilation amplitudes, which is currently poorly known and a source of significant uncertainty in many calculations. The present search is sensitive to both modes. Two–body charmless decays are also expected from bottom baryons. The modes $\Lambda_b^0 \rightarrow pK^-$ and $\Lambda_b^0 \rightarrow p\pi^-$ are predicted to have measurable branching fractions, of order $10^{-6}$, and, in addition to the interest in their observation, must be considered as a possible background to the rare $B_s^0$ and $B^0$ modes being investigated.

In this Letter we report the results of a search for rare decays of neutral bottom hadrons into a pair of charged charmless hadrons ($p, K$ or $\pi$), performed in 1 fb$^{-1}$ of $\bar{p}p$ collisions at $\sqrt{s} = 1.96$ TeV, collected by the upgraded Collider Detector (CDF II) at the Fermilab Tevatron. We report the first observation of modes $B_s^0 \rightarrow K^- \pi^+$, $\Lambda_b^0 \rightarrow pK^-$, and $\Lambda_b^0 \rightarrow p\pi^-$, and measure their relative branching fractions.

CDF II is a multipurpose magnetic spectrometer surrounded by calorimeters and muon detectors. The detector components relevant for this analysis are briefly outlined below: a more detailed description can be found in Ref. [11]. A silicon microstrip vertex detector (SVX) and a cylindrical drift chamber (COT) immersed in a 1.4 T axial magnetic field allow reconstruction of charged–particle trajectories (tracks) in the pseudorapidity range $|\eta| < 1.0$. The SVX consists of six concentric layers of double-sided silicon sensors with radii between 2.5 and 22 cm, each providing a measurement with up to 15 (70) $\mu$m resolution in the $\phi$ ($z$) direction. The COT has 96 measurement layers, between 40 and 137 cm in radius, organized into alternating axial and $\pm 2^\circ$ stereo superlayers. The transverse momentum resolution is $\sigma_{p_T}/p_T^2 \sim 0.15%/$(GeV/$c$), corresponding to a typical mass resolution of 22 MeV/$c^2$ for our signals. The specific ionization energy loss ($dE/dx$) of charged particles in the COT can be measured from the collected charge, which is logarithmically encoded in the output pulse width of each wire, and provides 1.5$\sigma$ separation between kaons and pions with momenta greater than 2 GeV/$c$.

The data were collected by a three-level trigger system, using a set of requirements specifically aimed at selecting two-pronged $B$ decays. At level 1, COT tracks are reconstructed in the transverse plane by a hardware processor (XFT) [13]. Two opposite-charge particles are required, with reconstructed transverse momenta $p_{T1}, p_{T2} > 2$ GeV/$c$, the scalar sum $p_{T1} + p_{T2} > 5.5$ GeV/$c$, and an azimuthal opening-angle $\Delta \phi < 135^\circ$. At level 2, the silicon vertex trigger (SVT) [14] combines XFT tracks with SVX hits to measure the impact parameter $d$ (distance of closest approach to the beam line) of each track with 45 $\mu$m resolution. The requirement of two tracks with 0.1 < $d$ < 1.0 mm reduces the light quark background by two orders of magnitude while preserving about half of the signal. A tighter opening-angle requirement, $20^\circ < \Delta \phi < 135^\circ$, preferentially selects two–body $B$ decays over multi–body decays with 97% efficiency and further reduces background. Each track pair is then used to form a $B$ candidate, which is required to have an impact parameter $d_B < 140$ $\mu$m and to have travelled a dis-
tance $L_T > 200 \mu m$ in the transverse plane. At level 3, an array of computers confirms the selection with a full event reconstruction. The overall acceptance of the trigger selection is $\approx 2\%$ for $b$-hadrons with $p_T > 4$ GeV/$c$ and $|\eta| < 1$.

The offline selection is based on a more accurate determination of the same quantities used in the trigger, with the addition of two further observables: the isolation ($I_B$) of the $B$ candidate [15], and the quality of the three-dimensional fit ($\chi^2$ with 1 d.o.f.) of the decay vertex of the $B$ candidate. Requiring a large value of $I_B$ reduces the background from light-quark jets, and a low $\chi^2$ reduces the background from decays of different long-lived particles within the event, owing to the good resolution of the SVX detector in the $z$ direction. The selection is optimized for detection of the $B_d^0 \to K^-\pi^+$ mode. Maximal sensitivity for both discovery and limit setting is achieved with a single choice of selection requirements [16] by minimizing the variance of the estimate of the branching fraction in the absence of signal [17]. The variance is evaluated by performing the full measurement procedure on simulated samples containing background and all signals from the known modes, but no $B_d^0 \to K^-\pi^+$ signal. The background fraction for each selection is determined from data by extrapolating the mass sidebands of the signal, and the signal yield is predicted by a detailed detector simulation. This procedure yields the final selection: $I_B > 0.525$, $\chi^2 < 5$, $d > 120 \mu m$, $d_B < 60 \mu m$, and $L_T > 350 \mu m$.

No more than one $B$ candidate per event is found after this selection, and a mass ($m_{\pi\pi}$) is assigned to each, using a charged pion mass assignment for both decay products. The resulting mass distribution is shown in Fig. 1. A large peak is visible, dominated by the overlapping contributions of the $B^0 \to K^+\pi^-$, $B^0 \to \pi^+\pi^+$, and $B^0 \to K^+K^-$ modes. A $B^0 \to K^+K^-$ signal would appear as an enhancement around 5.18 GeV/$c^2$, while signals for the other modes of this search are expected at masses higher than the main peak (5.33–5.55 GeV/$c^2$). Backgrounds include mis-reconstructed multi-body $b$–hadron decays (physics background) and random pairs of charged particles (combinatorial background).

We used an unbinned likelihood fit, incorporating kinematic (kin) and particle identification (PID) information, to determine the fraction of each individual mode in our sample. The likelihood for the $i$th event is

$$L_i = (1 - b) \sum_j f_j L_j^{\text{kin}} L_j^{\text{PID}} + b \left( f_p L_p^{\text{kin}} L_p^{\text{PID}} + (1 - f_p) L_c^{\text{kin}} L_c^{\text{PID}} \right),$$

where the index $j$ runs over all signal modes, and the index ‘p’ (‘c’) labels the physics (combinatorial) background terms. The $f_j$ are the signal fractions to be determined by the fit, together with the background fraction parameters $b$ and $f_p$.

The kinematic information is summarized by three loosely correlated observables: (a) the mass $m_{\pi\pi}$; (b) the signed momentum imbalance $\alpha = (1 - p_1/p_2)q_1$, where $p_1$ ($p_2$) is the lower (higher) of the particle momenta, and $q_1$ is the sign of the charge of the particle of momentum $p_1$; (c) the scalar sum of particle momenta $p_{tot} = p_1 + p_2$. The above variables allow evaluation of the invariant mass $m_{12}$ of a candidate for any mass assignment of the decay products $(m_1,m_2)$, using the equation

$$m_{12}^2 = m_{\pi\pi}^2 - 2m_{\pi}^2 + m_1^2 + m_2^2 + 2\sqrt{m_1^2 + m_2^2} \sqrt{p_1^2 + p_2^2} - 2\sqrt{p_1^2 + m_2^2} \sqrt{p_2^2 + m_1^2} - 2\sqrt{p_1^2 + p_2^2}.$$  \hspace{2cm} (2)

where $p_1 = \frac{1-|\alpha|}{2-|\alpha|}p_{tot}$, $p_2 = \frac{1}{2-|\alpha|}p_{tot}$.

We used the mass sidebands in data ($m_{\pi\pi} \in [5.00, 5.12] \cup [5.6, 6.2]$ GeV/$c^2$) to obtain the kinematic distributions of backgrounds [17]. The mass distribution of the combinatorial background is parameterized by an exponential function, while the physics background is modeled by an ARGUS function [18] convoluted with a Gaussian resolution function. In order to ensure the reliability of the search for small signals in the vicinity of larger peaks, the shapes of the mass distributions assigned to each signal have been modeled in detail. We have included the momentum dependence and non-Gaussian tails of resolution from a full simulation of the detector, and the effects of soft photon radiation in the final state, based on recent QED calculations [19]. This resolution model was checked against the observed shape of the $D^0 \to K^-\pi^+$ signal in a sample of $1.5 \times 10^6 D^{*+} \to D^0\pi^+$ decays, collected with a similar trigger selection. The observed discrepancies are below the $10^{-3}$.
mined from a sample of about 124,000 $\Lambda^0$ we measure all branching fractions relative to the production cross sections and absolute reconstruction efficiency, to allations of significantly smaller branching fractions. The model of the background allows for pion, kaon, proton, and electron components, whose fractions are determined from a sample of about 124,000 $\Lambda^0 \rightarrow p\pi^-$ decays. The model of the background allows for pion, kaon, proton, and electron components, whose fractions are determined by the fit. Muons are indistinguishable from pions with the available 10% fractional $dE/dx$ resolution and are therefore incorporated into the pion component.

From the signal fractions returned by the likelihood fit we calculate the signal yields shown in Table I. The significance of each signal is evaluated as the ratio of the yield observed in data, and its total uncertainty (statistical and systematic) as determined from a simulation where the size of that signal is set to zero. This evaluation assumes a Gaussian distribution of yield estimates, supported by the results obtained from repeated fits to simulated samples. This procedure yields a more accurate measure of significance with respect to the purely statistical estimate obtained from $\sqrt{-2\Delta \ln(L)}$. We obtain significant signals for the $B^0 \rightarrow K^-\pi^+$ mode (8.2$\sigma$), and for the $\Lambda^0_b \rightarrow p\pi^-(6.0\sigma)$ and $\Lambda^0_b \rightarrow pK^-(11.5\sigma)$ modes. Figure 2 shows relative likelihood distributions for these modes. No evidence is found for the modes $B^0 \rightarrow \pi^+\pi^-$ or $B^0 \rightarrow K^+K^-$, in agreement with expectations of significantly smaller branching fractions.

To avoid large uncertainties associated with production cross sections and absolute reconstruction efficiency, we measure all branching fractions relative to the $B^0 \rightarrow K^+\pi^-$ mode. Frequentist upper limits at the 90% C.L. are quoted for the unseen modes. For the measurement of $\Lambda^0_b$ branching fractions, the additional requirement $p_T(\Lambda^0_b) > 6$ GeV/c was applied to allow easy comparison with other $\Lambda^0_b$ measurements at the Tevatron, which are only available above this threshold $[20,23]$. This additional requirement lowers the $\Lambda^0_b$ yields by about 20%. The raw fractions returned by the fit were corrected for the differences in selection efficiencies between different modes, which range from 8% to 40% for the measurements of $b$-mesons and $\Lambda^0_b$ branching fractions, respectively. These corrections were determined from detailed detector simulation, with the following exceptions that were measured from data: the momentum-averaged relative isolation efficiency between $B^0_1$ and $B^0_2$, 1.00 $\pm$ 0.03, has been determined from fully-reconstructed samples of $B^0 \rightarrow J/\psi\phi$, and $B^0 \rightarrow J/\psi K^{*0}$ decays $[17]$. The difference in efficiency for triggering on kaons and pions due to the different specific ionization in the COT (a $\approx 5\%$ effect) was measured from a sample of $D^+ \rightarrow K^-\pi^+\pi^+$ decays triggered on two tracks, using the unbiased third track $[24]$. Possible differences in efficiency of the isolation requirement between $B^0$ and $\Lambda^0_b$, and in the trigger efficiency between kaons and protons, were taken into account in the systematic uncertainties.

The dominant contributions to the systematic uncertainty are the uncertainty on the combinatorial background model and the uncertainty on the $dE/dx$ cali-

![Figure 2: Distribution of the relative signal likelihood, $\mathcal{L}_{S}/(\mathcal{L}_{S} + \mathcal{L}_{\text{other}})$, in the region $5.1 < m_{\pi\pi} < 5.6$ GeV/c$^2$. For each event, $\mathcal{L}_{S}$ is the likelihood for the $B^0 \rightarrow K^-\pi^+$ (a), $\Lambda^0_b \rightarrow pK^-$ (b), or $\Lambda^0_b \rightarrow p\pi^-$ (c) signal hypotheses, and $\mathcal{L}_{\text{other}}$ is the likelihood for everything but the chosen signal, i.e. the weighted combination of all other components according to their measured fractions. Points with error bars show the distributions of data and histograms show the distributions predicted from the measured fractions.]

<table>
<thead>
<tr>
<th>Mode</th>
<th>$N_s$</th>
<th>Significance</th>
</tr>
</thead>
<tbody>
<tr>
<td>$B^0_1 \rightarrow K^-\pi^+$</td>
<td>230 $\pm$ 34 $\pm$ 16</td>
<td>8.2$\sigma$</td>
</tr>
<tr>
<td>$B^0_2 \rightarrow \pi^+\pi^-$</td>
<td>26 $\pm$ 16 $\pm$ 14</td>
<td>$&lt; 3\sigma$</td>
</tr>
<tr>
<td>$B^0_2 \rightarrow K^+K^-$</td>
<td>61 $\pm$ 25 $\pm$ 35</td>
<td>$&lt; 3\sigma$</td>
</tr>
<tr>
<td>$\Lambda^0_b \rightarrow pK^-$</td>
<td>156 $\pm$ 20 $\pm$ 11</td>
<td>11.5$\sigma$</td>
</tr>
<tr>
<td>$\Lambda^0_b \rightarrow p\pi^-$</td>
<td>110 $\pm$ 18 $\pm$ 16</td>
<td>6.0$\sigma$</td>
</tr>
</tbody>
</table>
bration and parameterization. Other contributions come from trigger efficiencies, physics background shape and kinematics, $b$–hadron masses and lifetimes, and the possible polarization of $Λ_b^0$ decays.

The final results are listed in Table II. Absolute branching fractions are also quoted, by normalizing to world-average values of production fractions and $B(B^0 \to K^+ \pi^-)$ [1, 21]. The branching fraction measured for the $B^0_s \to K^- \pi^+$ mode is consistent with the previous upper limit (< $5.6 \times 10^{-6}$ at 90% C.L.), based on a subsample of the current data. This agrees with the prediction in Ref. [27], but it is lower than most other predictions [5, 6, 24]. The $B^0_s \to \pi^+ \pi^-$ upper limit improves and supersedes the previous limit [3]. The present measurement of $B(B^0 \to K^+ K^-)$ is in agreement with other existing measurements and has a similar resolution [21], but the resulting upper limit is weaker due to the observed central value. The sensitivity to both $B^0 \to K^+ K^-$ and $B^0_s \to \pi^+ \pi^-$ is now close to the upper end of the theoretically expected range [3, 6, 7, 27]. We also report the first branching fraction measurements of charmless $Λ_b^0$ decays. They are significantly lower than the previous upper limit of $2.3 \times 10^{-5}$ [28], and in reasonable agreement with predictions [3], thus excluding the possibility of large ($O(10^6)$) enhancements from R-parity violating supersymmetric scenarios [29]. Their ratio can be determined directly from our data with greater accuracy than the individual values. For this purpose, the additional $p_T > 6$ GeV/c is not necessary, and we can exploit the full sample size, obtaining $B(Λ_b^0 \to p \pi^-)/B(Λ_b^0 \to p K^-) = 0.66 \pm 0.14 \pm 0.08$, in good agreement with the predicted range 0.60–0.62 [1].

In summary, we have searched for rare charmless decay modes of neutral $b$–hadrons into pairs of charged hadrons in CDF data. We report the first observation of the modes $B^0 \to K^- \pi^+$, $Λ_b^0 \to p \pi^-$, and $Λ_b^0 \to p K^-$, and measure their relative branching fractions. We set upper limits on the unobserved modes $B^0 \to K^+ K^-$ and $B^0_s \to \pi^+ \pi^-.$

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<table>
<thead>
<tr>
<th>Mode</th>
<th>Relative $B$</th>
<th>Absolute $B$ (10$^{-6}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>$B^0 \to K^+ \pi^- ,$</td>
<td>$= 0.071 \pm 0.010 \pm 0.007$</td>
<td>$5.0 \pm 0.7 \pm 0.8$</td>
</tr>
<tr>
<td>$B^0 \to \pi^+ \pi^- ,$</td>
<td>$= 0.007 \pm 0.004 \pm 0.005$</td>
<td>$0.49 \pm 0.28 \pm 0.36$ (1.2 at 90% C.L.)</td>
</tr>
<tr>
<td>$B^0 \to K^- K^- ,$</td>
<td>$= 0.020 \pm 0.008 \pm 0.006$</td>
<td>$0.39 \pm 0.16 \pm 0.12$ (0.7 at 90% C.L.)</td>
</tr>
<tr>
<td>$Λ_b^0 \to p K^- ,$</td>
<td>$= 0.066 \pm 0.009 \pm 0.008$</td>
<td>$5.6 \pm 0.8 \pm 1.5$</td>
</tr>
<tr>
<td>$Λ_b^0 \to p \pi^- ,$</td>
<td>$= 0.042 \pm 0.007 \pm 0.006$</td>
<td>$3.5 \pm 0.6 \pm 0.9$</td>
</tr>
</tbody>
</table>

[12] CDF II uses a cylindrical coordinate system in which $\phi$ is the azimuthal angle, $r$ is the radius from the nominal beam line, and $z$ points in the proton beam direction, with the origin at the center of the detector. The transverse plane is the plane perpendicular to the $z$ axis.


[15] Isolation is defined as $I_B = p_T(B)/(p_T(B) + \sum p_T_i)$, where $p_T(B)$ is the transverse momentum of the $B$ candidate, and the sum runs over all other tracks within a cone of radius 1, in $\eta$-$\phi$ space around the $B$ flight-direction.


