

1 Acquisition and evaluation of radiometrically comparable multi-footprint airborne LiDAR data for forest remote sensing

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3  
4 Abstract

5 Forest inventories comprise observations, models and sampling. Airborne LiDAR has established its role in providing  
6 observations of canopy geometry and topography. These data are input for estimation of important forest properties to support  
7 forestry-related decision-making. The primary deficiency in forest remote sensing is tree species identification. This study  
8 examines the quite atypical option of using multi-footprint airborne LiDAR data. Features of such sensor design exist in recently  
9 introduced multispectral laser scanners. The first objective was to acquire radiometrically normalized, multi-footprint (11, 22, 44  
10 and 59 cm) waveform (WF) data that characterize 1064-nm backscatter reflectance on the interval scale. The second objective  
11 was to analyze and validate the data quality in order to draw the correct conclusions about the effect of footprint size. Finally,  
12 the data were analyzed in different forest canopies. The experiment was carried out in Finland. Footprint variation was  
13 generated by acquiring data at different flying heights and by adjusting the transmitted power. The LiDAR campaign was  
14 successful and the data were of sufficient quality, except for a 1 dB trend due to the atmosphere. Significant findings were made  
15 concerning the magnitude of atmospheric losses, the linearity of the amplitude scale and the bandwidth characteristics of the  
16 receiver, the stability of the transmitter, the precision of the amplitude data, the transmission losses in canopies and power lines  
17 as well as the response of WF attributes to footprint size in forest canopies. Multi-footprint data is a promising approach  
18 although the species-specific signatures were weak.

19 Keywords: Canopy structure, Tree species, Forest inventory, Laser scanning, Waveform, Normalization, Linearity, Impulse  
20 response, Losses

## 21 1. Introduction

22  
23 Airborne LiDAR data are used for many purposes, and the rationales for the present study originate from LiDAR remote sensing  
24 (RS) of forests (Vauhkonen et al., 2014). It has recently become an integral part of forest management planning systems in  
25 Finland (Maltamo and Packalén, 2014). LiDAR is an observation tool that has reduced the sampling intensity and provided  
26 entirely new observations for the estimation of forest attributes. In Finland, forest planning inventories consist of area-based  
27 estimation of stand boundaries, attributes and management proposals, using a combination of field reference, aerial image and  
28 LiDAR data, all finally followed by a field inspection. The data acquisition costs have decreased compared with field-work  
29 intensive systems, with both improvements and deficiencies in the deliverables. The primary deficiency in forest RS, not only in

30 Finland, is tree species identification. There is a need for observations that aid in solving this problem. Multi-footprint data may  
31 comprise one option.

32 Regarding sensors, foresters have mostly used data acquired by topographic small-footprint instruments. Narrow beams  
33 promote accurate geometry and reach the ground. While simulation studies have highlighted how sensor properties influence  
34 the available information (Disney et al., 2010, Hovi & Korpela, 2014), certain constraints pertain to sensor design. They include  
35 the need for extremely high dynamic range, consideration of eye-safety, receiver sensitivity and bandwidth, availability of  
36 powerful lasers at different wavelengths, data transfer and storage capacity requirements, power consumption, limitations on  
37 the size of the mirror and the aperture, and the final instrument cost. Advances in the last 10–15 years relate to improvements  
38 in pulse repetition frequency, capture of waveform (WF) data, fast direct georeferencing and the detection of multiple echoes,  
39 inter alia. Teledyne Optech (Vaughan, Ontario) recently released a three-band, dual-footprint WF-recording sensor called Titan,  
40 in which the divergence of the 532-nm pulses is wider than that of 1064-nm and 1550-nm pulses. The relative reflectance  
41 difference of green and woody vegetation between the 1064 and 1550 nm bands is likely beneficial for tree species iden-  
42 tification as the green/woody silhouette visible from the above shows between-species differences owing to differences in leaf  
43 angle distributions, and in branch and crown morphology. On the other hand, WFs carry species-specific traits (Hovi et al., 2015),  
44 when tens of pulses per crown are evaluated. How the combination of WF and multi-band data will benefit species  
45 identification, is a very interesting topic. Although advances are foreseeable with multispectral LiDAR, this study examines the  
46 relatively unconventional concept of multi-footprint LiDAR to investigate how such data characterize canopy structure, and in  
47 general augment our understanding of the characteristics of LiDAR data. The rest of this section provides further background to  
48 the topic and introduces hypotheses that were tested empirically.

49 The radar equation (Eq. 1) applies to LiDAR as well, and states that the received power ( $P_r$ ) is influenced by the aperture area ( $D$ ),  
50 the atmospheric transmission losses ( $T$ ), the optical efficiency of the system ( $Q$ ), the transmitted power ( $P_t$ ), the laser beam  
51 divergence ( $\beta$ ), the range ( $R$ ), and the backscattering cross-section of the target ( $\sigma$ ):

$$52 \quad P_r = \frac{P_t D Q}{4\pi\beta^2} T^2 \frac{\sigma}{R^4} \quad (1),$$

53 where  $\sigma$  can be said to be comprised of target reflectance, geometry and the illuminated/silhouette area. Eq. (1) is  
54 'instantaneous', while in reality the power terms and  $\sigma$  are time- or distance dependent and the optical power of the beam has  
55 a Gaussian cross-section profile (cf. Wagner et al., 2006; Mallet and Bretar, 2009).

56 Eq. (1) states that  $P_r$  is in linear dependence with target reflectance and the transmitted power. A large aperture increases  $P_r$  and  
57 the signal-to-noise ratio (SNR), because the extra solar illumination reaching the detector, non-attenuated, is minor compared to  
58 the internal receiver noise. Atmospheric attenuation is 0.1–0.3 dB/km in weather conditions that are suitable for LiDAR RS.

59 The effect of  $R$ , i.e. the spherical losses, depend on the target, because large surfaces, linear targets and blob-like objects all  
60 exhibit different response. A sensor that provides ratio- or interval-scale measurements of  $P_r$ , will show a fourfold (6 dB)  
61 increase in well-defined surfaces, when  $R$  is halved, whereas  $P_r$  is 8 or 16 times (9 or 12 dB) larger for a wire or a small leaf. Tree  
62 crowns are volumetric and the structure influences the signal change due to  $R$  (Korpela et al., 2010, Gatzliolis, 2011), making  
63 range normalization ill-posed in canopies. In linear targets, such as cables, the distribution of  $P_r$  has an exponential-even shape  
64 with a modest peak near the maximum, while for blob-like features, the distribution is inverse-exponential.

65 Divergence defines the angular spread of the beam's Gaussian irradiance profile. Here it is defined by the  $1/e^2$  point. A target  
66 may give rise to a measurable echo, only if it is illuminated by the pulse center. The backscatter cross-section may shrink for a  
67 tilted target, or at least the returning pulse is extended and dampened. Because sensors record meaningful data only,  
68 backscattering ( $P_r$ ) needs to be strong enough to trigger an observation, and the real signal distributions are always truncated by  
69 some combination of  $\sigma$  and the receiver noise. The sensor used in this study stores a single continuous and fixed-length WF, the  
70 recording of which is triggered by the first threshold-exceeding echo. Some sensors record variable-length WF samples with  
71 pauses in between and thresholding applies to this 'piecewise WF-recording' as well (e.g. Armston et al., 2010). In general, weak  
72 backscattering preceding the first echo may remain undetected, while such weak within-canopy or ground scattering may be  
73 observable, when the recording is on, and missed, while it is paused. This constitutes a vertical sampling bias. A somewhat  
74 similar effect causes selection bias in the analyses of discrete-return (DR) intensity data, when the separation between only and  
75 first-of-many echoes is often made (Ørka et al., 2010; Korpela et al., 2010b), because this instrument-driven division is  
76 influenced by the backscatter cross-section profile of targets following the first interaction, for example the ground flora in the  
77 case of tree crowns.

78 Regarding multi-footprint data, changing solely the beam width will not influence  $P_r$  as long as the target is a well-defined  
79 surface filling the footprint, while for small targets,  $P_r$  varies according to the intersection geometry, as stated. Multiple  
80 scattering and transmission losses are introduced instantly following the first pulse-canopy interaction (the subsequent targets  
81 comprise an apparent differential backscatter cross-section profile (Richter et al., 2015), high-order volumetric scattering can be  
82 significant, and backscattering occurs far from the beam path (Hovi and Korpela, 2014). Foliage orientation, size, density and  
83 silhouette area were shown to influence the recorded intensity in broadleaved and coniferous trees in Finland (Korpela et al.,  
84 2010b, 2013; Hovi and Korpela, 2014). If beam width increases by a factor of  $i$ , while  $P_t$  is constant, the average  $P_r$  decreases by  
85  $1/i^2$  and  $1/i$  in blob-like and linear targets, respectively. A wide beam is more likely to find directional canopy gaps and the  
86 proportion of pulses that illuminate the ground is higher. Many small targets that trigger an echo in narrow beams remain  
87 undetected in large-footprint data and contribute to transmission losses, which explain why ground signal levels can be  
88 expected to be lower in large-footprint data. The wider the pulse, the smaller is the likelihood for the pulse to reach the ground  
89 with maximal energy. Thus, measuring the reflectance properties of the forest floor using single-return data is more accurate in

90 narrow footprint data as shown for understory lichens in [Korpela \(2008\)](#). The peak amplitude distribution of canopy echoes can  
91 therefore be expected to be shifted towards low values in large footprint data. It can also be expected that the return WFs are  
92 widened in large-footprint data, because a wider beam illuminates targets that also spread across a larger depth.

93 As stated, spherical losses depend on the target geometry and it is therefore evident that a fixed signal-to-noise-ratio (SNR)  
94 cannot be maintained for the canopy and the ground, when scanning from different heights ([Goodwin et al., 2006](#), [Ørka et al.,  
95 2010](#)). When  $R$  is doubled, a 6 dB increase in  $P_t$  will maintain the SNR for well-defined surfaces. We cannot expect this to hold  
96 true for the forest floor, because canopy transmission losses will increase with increasing footprint size. The SNR of small targets  
97 will fall even if the power is increased by 6 dB, which explains the reduction in canopy echo counts in high-altitude LiDAR  
98 ([Goodwin et al., 2006](#)). Of course, atmospheric losses introduce additional signal loss. Varying the scan zenith angle changes  $R$   
99 and the incidence angles. Pulses reflecting from tilted surfaces will exhibit widened echoes, in which the peak amplitude is dam-  
100 pened, and both the widening and the dampening depend on the footprint size (e.g. [Mallet and Bretar, 2009](#)).

101 Increasing  $P_t$  improves the SNR. Other options include enlarging the receiver aperture and mirror, or decreasing the receiver  
102 noise level. The latter comprises dark noise of the detector, electronic and digitization noise as well as natural light that reaches  
103 the detector. The energy captured by an aperture that is located in the footprint center, and focuses light to the retina, cannot  
104 exceed critical values, which in turn limits  $P_t$  and/or beam divergence. Eye-safety is a more severe concern at 532 nm than at  
105 1064 nm or 1550 nm. Regarding multi-band sensors, we can note that the reflectance of foliage at 532 nm is one tenth of that at  
106 1064 nm. Given the same aperture, divergence and output power, the green band signals will thus be 10 times lower. In order to  
107 improve the SNR of well-defined targets (at 532 nm), it is possible to increase both the transmitted power and/or beam  
108 divergence, where the latter may be needed for eye-safety. Because multiple scattering in vegetation is weak in visible bands,  
109 the returning 532 nm pulses are probably less extended compared to NIR pulses of the same beam divergence, which may be  
110 exploited in target classification using multi-band LiDAR.

111 Radiometric observables in LiDAR data include intensity and WF amplitude data. The instantaneous spectral irradiance is  
112 captured by the aperture, reflected to a collimating lens that directs it to the photon detector through a band-pass filter. The  
113 detector and the circuits that follow, i.e. the receiver, have certain noise characteristics, impulse response and bandwidth (BW)  
114 properties. The latter two determine how the dynamics of the photon surge are captured. Receiver bandwidth is inherently  
115 limited and rapid changes in the photon surge are low-pass filtered in the WF. As long as the transmitted pulses are stable,  
116 minor deficits in bandwidth are not critical for reliable ranging, which is also shown in this study. Receiver performance may  
117 depend on signal strength and for example very strong echoes can temporarily increase the noise level or cause 'ringing', i.e.  
118 'ghost echoes', in the WF (e.g. [Armston et al., 2010](#)). If the internal signal levels are too high, i.e., outside the linear (dynamic)  
119 range of the receiver circuits, the sensor of this study performs erroneous discrete ranging and the intensity data are distorted.

120 A one-nanosecond sampling interval is typical in WF-recording sensors, and the length (full width at half maximum, FWHM) of  
121 the transmitted pulse is below 10 ns. The power goes up to several kW in some nanoseconds only, but the exact power  
122 envelope of the transmitted pulse is usually not known. Any within- or between-pulse variation directly influences the  
123 convolution of the transmitted pulse with the backscatter cross-section profile of the targets. Some sensors provide a digitized  
124 sample of the transmitted pulse to enable exact ranging and echo attribute retrieval in post processing (Wagner et al., 2006;  
125 Roncat et al., 2014)

126 Radiometrically quantitative analyses require that the properties of the transmitted pulse (directional target illumination) and  
127 the receiver's response to instantaneous at-sensor irradiance (measurement of reflected radiance) are known, in addition to the  
128 properties of the medium. Such absolute calibration is typically unattainable for LiDAR sensors and vicarious calibration with  
129 reflectance targets becomes the only option (cf. Wagner, 2010). This in turn is challenging because reference measurements of  
130 zero phase-angle reflectance are hard to establish and are rarely done. A linear response of the receiver is usually assumed (with  
131 gain and offset), or alternatively a look-up-table is used. Some sensors have two receivers and the model is needed for both. DR  
132 sensors can have individual circuits for the detection of echoes and the intensity values can have different calibration. The  
133 sensor used in this study has a single receiver, in which the gain of the signal chain is dynamically adjusted up to 3 dB by an  
134 automatic gain control (AGC) circuit. In order to draw conclusions about the influence of footprint size, using the WF data, it is  
135 vital to establish a mapping between the amplitude data and the at-sensor power. Such mapping was performed in this study.

136 This study examines the relatively unconventional concept of multi-footprint LiDAR to investigate how such data characterize ca-  
137 nopy structure, and in general augment our understanding of LiDAR data. Specific objectives were:

- 138 1. Plan and carry out WF LiDAR data acquisition with a fixed beam divergence sensor that results in radiometrically  
139 comparable multi-footprint data in well-defined targets.
- 140 2. Carry out radiometric normalization of peak amplitude data using well-defined surfaces to remove the effects of the  
141 AGC circuit, to estimate the offset of the amplitude scale and to identify any deficits in the acquisition.
- 142 3. Validate the normalization across the study site to ensure that the data qualifies for analyses on the influence of the  
143 footprint size.
- 144 4. Acquire independent data using the same sensor to investigate and verify whether the amplitude scale of the receiver  
145 is in linear dependence with the instantaneous at-sensor power.
- 146 5. Investigate the impulse response of the receiver to assure that WF attributes are influenced only by the pulse-target  
147 interaction.
- 148 6. Carry out empirical tests to show that the hypotheses presented in the introduction are confirmed by the measured  
149 real data.

## 2. Experiment and results

### 2.1 Research site

The experiments were carried out in Hyytiälä (61°50'N, 24°17'E), southern Finland. The area is mostly covered by forest, but there are also fields, open and forested peatlands, built environment and a network of forest roads. Aerial photographs cover the time period of 1946–2015 and there are 19 airborne LiDAR datasets acquired in 2004–2015. Geometric reference data include close-range photogrammetric images, surveying points and profiles. These data as well as the redundant airborne image and LiDAR data were used for geometric quality control of the LiDAR data sets used here. (Hovi, 2015).

### 2.2 Waveform-recording LiDAR data

Objective #1 was to acquire multi-footprint data, in which the same energy was transmitted by varying the beam divergence. Sensors with the option of changing divergence were not available. To produce the same effect, a Leica ALS60 sensor was used at different heights as its transmitted power can be adjusted by almost 20 dB (1–100%). The nadir  $1/e^2$  footprint diameters were 11, 22, 44 and 59 cm from the acquisition heights of 0.5, 1.0, 2.0 and 2.7 km, respectively. The pulse power was set to nominal levels of 100, 55, 14 and 4% (Tables 1 and 2) to account for the spherical losses in well-defined surfaces. Strip overlaps were increased for the higher altitudes to reach comparable pulse density. According to the manufacturer, the FWHM of the transmitted pulses is 4 or 9 nanoseconds, which showed as 7 and 10 ns in the recorded data (Fig. 1). The wavelength is 1064 nm, and Fig. 2 shows the difference of intensity data in different targets between the ALS60 and a Riegl LMS-Q680i sensor, which operates at 1550 nm. The Riegl data was acquired in 2013 (not tabulated) and was used for illustrations only. The same ALS60 sensor was used for datasets 2012\_## and 2013\_## (Table 2). The acquisition in 2013 was carried out to verify the linear response of the receiver, objective #4. It comprised three flying heights: 700 m, 800 m and 900 m. The sensor was operated in an unorthodox manner by deliberately fixing both the transmitter power and the receiver gain (AGC, see later). Appropriate gain and power settings were found by first acquiring the 700-m data with the AGC on. These datasets have  $\pm 12.5\%$  range variation.

The receiver front-end in the ALS60 implements an AGC and the gain can vary by 0–3.2 dB (typically less) in response to scene brightness so that the internal signal levels are kept in the linear region of the range detection circuits. As said, the gain can also be fixed (2013\_## datasets). The influence of the AGC can largely be removed, using an 8-bit 'voltage value' that is stored for each pulse, shown first in Korpela (2008).

177 Table 1. Sensor parameters.

Wavelength, nm	1064
Pulse length, FWHM, ns	4, 9
System WF, FWHM, ns	7.8, 10.2
WF of transmitted pulse	No
Divergence, mrad, $1/e^2$	0.22
Mirror	Oscillating
Front-end amplifier	Automatic Gain Control
Triggering of WF-recording	Discrete-return circuit
WF-recording sequence	Single, 256 samples
WF Sampling rate, ns	1 or 2
LiDAR position, XYZ	yes
Discrete-return echoes	on-the-fly; 1, 2, 3 or 4
Echo attributes	range, XYZ, intensity
Acquisition height, m	500, 1000, 2000, 2700

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179

180 Table 2. Characteristics of the experimental LiDAR acquisitions. See also Table 1.

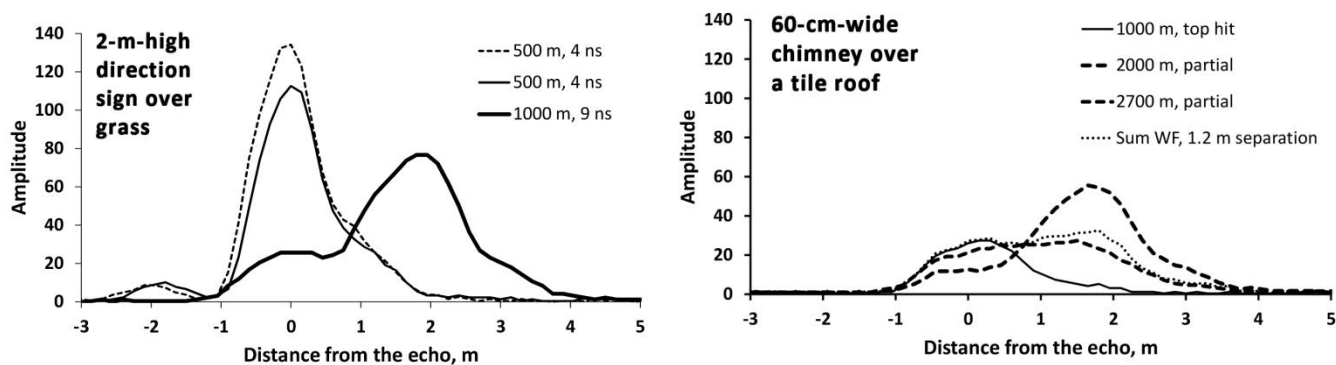
Dataset	2012_ (05, 10, 20, 27)	2013b_ (07, 08, 09)
Purpose	Multi-footprint data	Linearity tests
Date, local	July 5, 2012	June 15–16, 2013
Time, GMT	19–22	21–00
Phenology	mid-season	early season
Sensor	ALS60	ALS60
Altitude, m	500, 1000, 2000, 2700	700, 800, 900
WF density per $m^2$	5, 5, 3, 2	6, 6, 5
Scan angle, degrees	$\pm 15$	$\pm 15$
PRF, kHz	152, 99, 59, 45	106
System WF, FWHM, ns	7.8, 10.1, 10.3, 10.4	7.8
Footprint, $1/e^2$ , m	0.11, 0.22, 0.45, 0.59	0.15, 0.18, 0.20
AGC	on	off

181

182 ALS60 is primarily a DR sensor and the WF storage calls for an optional digitizing module. The storage is triggered by the first DR  
183 echo. A sequence of 256 amplitude (A) values is recorded including a buffer of 30 samples (Fig. 1). Two oscilloscopes take turns  
184 in digitizing the pulses at 1 GS/s. Their bandwidth (BW) is 300 MHz and the rise time for a square signal from 10% to 90% is 1.2  
185 ns. The BW and linearity of the preceding circuits, including the impulse response of the photon detector are unknown, as is the

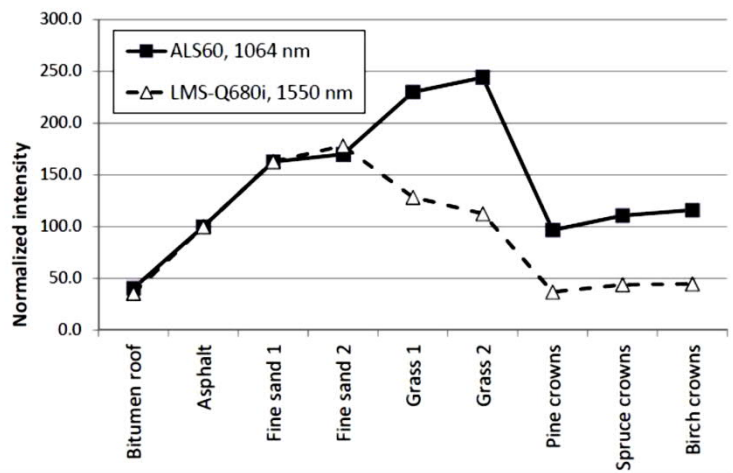
186 exact stability of the transmitter. However, the large difference in the FWHM of the transmitted and received 4 ns pulses clearly  
 187 suggests bandwidth deficits.

188



189

190 Fig. 1. The x-axis shows the range in meters centered (0 m) at the first echo. In the left figure there are three single-echo WFs.  
 191 The 500-m echoes are from the ground, while the 1000-m echo was triggered 2 m above the ground despite the strong ground  
 192 signal. Weak backscattering from the signs is seen even for the 500-m pulses in the preceding buffer. In the figure on the right  
 193 there are three pulses that gave rise to an echo from the top of the chimney. The 1000-m pulse reflected 'entirely' from the top,  
 194 while the 2/2.7 km pulses have scattered from the chimney and the roof. The 'Sum WF' is the 1000-m WF summed with its  
 195 shifted (6 ns, 90 cm) copy.



196

197 Fig. 2. Averaged and normalized peak amplitude values (asphalt=100) in 1064 and 1550 nm pulses (not tabulated, Riegl LMS-  
 198 Q680i sensor, 700 m AGL, amplitude data linearized using a look-up-table) in different surfaces. Major reflectance differences of  
 199 green vegetation are seen between 1064 and 1550 nm. Backscattering from tree crowns is not only influenced by the  
 200 reflectance properties but also by footprint size and the foliage orientation and crown structure. Footprint diameters here are  
 201 comparable, 0.44 m and 0.40 m.

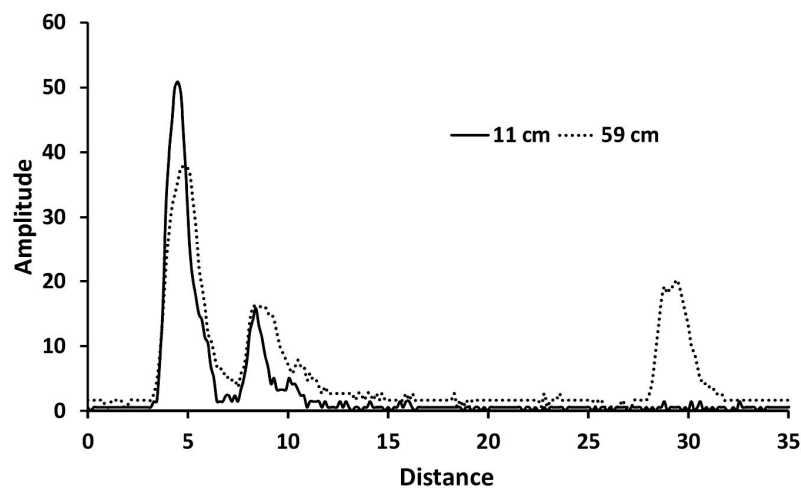
202 [2.3 Waveform attributes](#)



203 The WF attributes were computed in a straightforward way for what were called noise-exceeding amplitude sequences (NEAS,  
204 also referred to as an echo). More advanced methods for detecting individual echoes that may contribute to a NEAS, and for the  
205 derivation of echo attributes such as amplitude and width are discussed in e.g. Roncat et al. (2014). A detailed description of the  
206 calculations is given in Hovi et al. (2015). In Hovi et al. (2015) these were found important in tree species identification.

207 Noise was defined as the variation around a baseline signal, during which the receiver recorded background illumination only.  
208 The AGC was a particular challenge in defining the baseline, as the gain influenced both the baseline and the noise. The baseline  
209 varied by up to 0.3 amplitude units, which is only 0.25% of the maximum amplitude and neglecting it did not influence the  
210 results. The threshold for the start and end of a NEAS was three standard deviations of the observed variation around the  
211 baseline (Fig. 3).

212 In brief, the WF is first moderately low-pass filtered. The NEAS is found by tracing the (interpolated) points where the signal  
213 arises from and falls into the noise. The highest amplitude in a NEAS is the peak amplitude,  $pA$ , which is in strong dependence  
214 with discrete-return intensity, but only in well-defined targets. The sum of the amplitude values is energy,  $E$ . The width  
215 attributes were FWHM and  $L$ . FWHM was computed for the highest peak (defining the half maximum) and  $L$  is the total length of  
216 a NEAS. The rise speed,  $riS$ , was assumed to describe 'the porosity' of the illuminated target, and was defined as the ratio of the  
217 rise time and that of a pulse reflecting from a well-defined surface having the same  $pA$ . The normalization is needed as the  
218 above-the-noise rise time depends on the signal strength (c.f. constant fraction discriminator in range detection).



219  
220 Fig. 3. WFs of two adjacent and collinear pulses in a pine forest. The 59-cm WF is shifted up by one unit to be visible. The first  
221 echoes have reflected from a branch 24.7 m high and had an XYZ separation of 0.16 m, while the ground intersection points had  
222 an XY separation of 0.17 m. The scan zenith angle was 8 degrees. Ground echo can be seen for the wider pulse. The attribute  
223 values for the first NEASs are 50 and 37 for  $pA$ , 8.3 and 12.6 for FWHM, 1.16 and 1.16 for  $riS$ , and 486 and 518 for  $E$ .

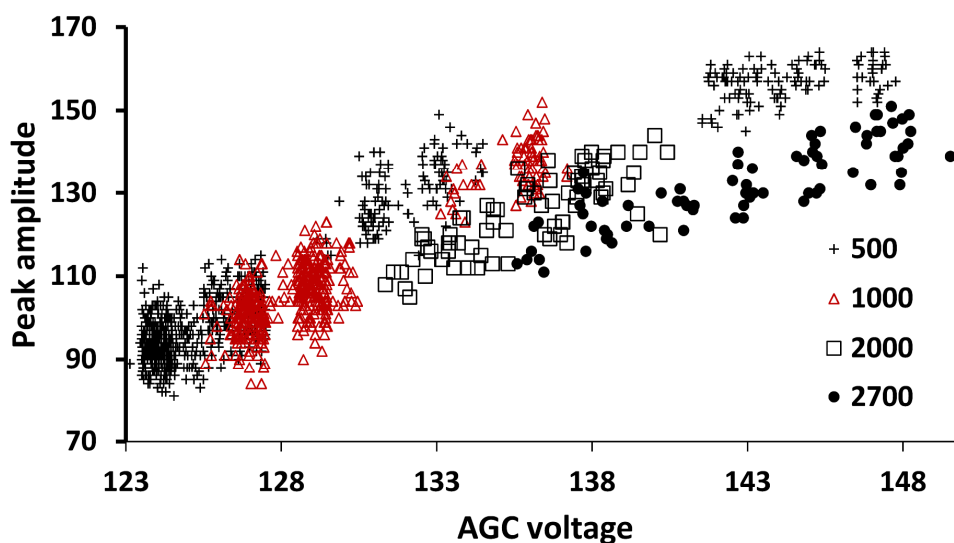
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## 2.4 Radiometric normalization of the multi-footprint data

227 As stated, the AGC amplifier is a particular property of the receiver in the ALS60. In the datasets used, the relative gain was  
 228 1–2.1, i.e. 0–3.2 dB in power, assuming that the intensity and amplitude values are in linear dependence with at-sensor power.  
 229 The influence of the AGC was the same on amplitude and intensity data. The  $AGC_{voltage} \times gain$  relationship resembles a sigmoid,  
 230 and the parameters of a normalization function ( $f_{AGC}$ , a 2nd-degree polynomial) were estimated using homogenous surfaces of  
 231 varying reflectance (Fig. 4). These surfaces were bitumen roof, old asphalt, fine sand and grass. Fig. 4 shows the uncorrected  
 232 data for grass.



233

234 Fig. 4. Uncorrected peak amplitude x  $AGC_{voltage}$  distributions of pulses reflecting from grass in the four 2012\_## datasets. The  
 235 sigmoid shape is seen in the 500 m data.

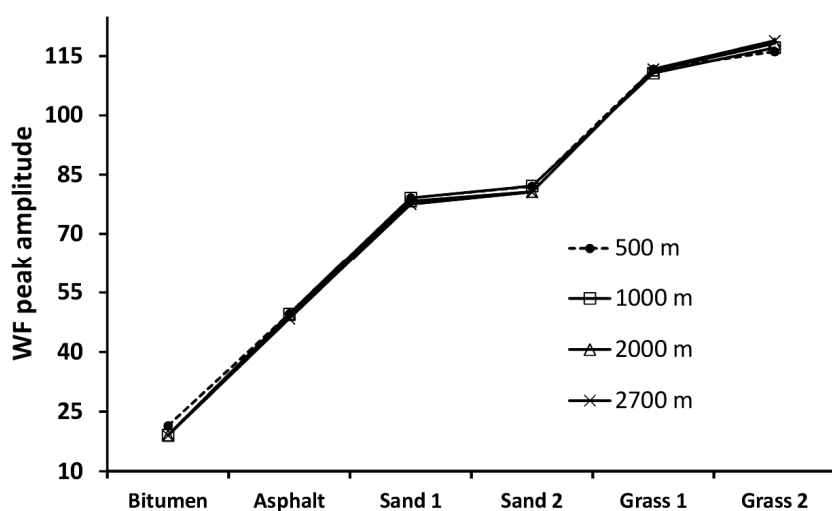
236 The normalizing model (Eq. 2) of a pulse  $i$ , acquisition height  $j$  and sensor  $k$ , used the explanatory variables  $R$ ,  $AGC_{voltage}$  and the  
 237 scan zenith angle. The parameters to be solved were the offset of the amplitude, the parameters of  $f_{AGC}$ , as well as coefficients  $d_j$   
 238 for each acquisition height that absorb possible deviations of the transmitted power, recalling that the power could only be set  
 239 at 1% nominal intervals.

$$240 \quad A_{norm(i,j,k)} = \left( \frac{R_i}{R_{ref(k)}} \right)^2 \cdot \frac{1}{f_{AGC(k)}} \cdot (A_{obs(i)} - offset_k) \cdot \cos(\theta_i) \cdot d_j \quad (2)$$

241 The estimation of  $f_{AGC}$  was carried out by first minimizing a weighted sum of coefficients of variation and trend coefficients of  
 242 univariate regression, indicating removal of trends, over all surfaces and datasets. At this phase 80–95% of the within-surface  
 243 variance was removed. Parameters  $offset$  and  $d_j$  were finally estimated by minimizing per surface mean amplitude differences of  
 244 different acquisition heights in each surface (Fig. 5). In the solution,  $d_j$  were 1.00, 3.56, 12.94, and 22.03 and the offset of the A

245 scale was 11.4. If the power had been set perfectly and atmospheric attenuation was zero, the expected values for  $d_j$  are 1, 4,  
246 16, and 29.16. Unfortunately, it was realized at this phase that the flight planning software did not account for atmospheric  
247 losses. Using the  $d_j$  values, attenuation was assessed to be  $0.22 \pm 0.06$  dB/km, which is realistic for 1064 nm. The 0.06 dB/km  
248 residual standard deviation shows that losses explain the deviations only partly. Power adjustment was done at 1% intervals,  
249 which corresponds to 25% relative accuracy in the 500 m data. Overall, the data showed an approximately 1 dB trend from 0.5  
250 to 2.7 km, which is also visible in Fig. 4 as 'vertical layers'.

251 Parameter estimation (vicarious calibration) of Eq. 2 could not be done by constraining with the absolute reflectance differences  
252 of the calibration surfaces, as only HCRF (hemispherical-conical reflectance factor) measurements at 900 nm by a field  
253 spectrometer were available. These measurements were shown to match the 2012 AGC-corrected data reasonably well in  
254 Korpela et al. (2013). However, the between-surface differences in Fig. 5 are in the instrument scale. Thus, the linearity of the A  
255 data, with respect to power entering the aperture, had to be tested (Section 2.5).



256

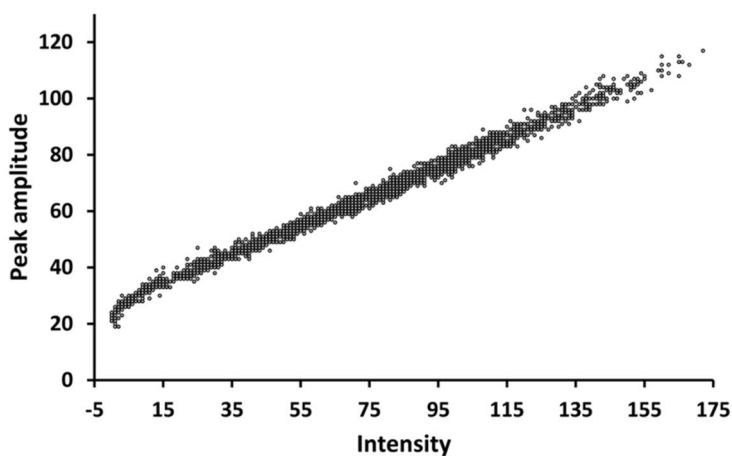
257 Fig. 5. Normalized mean (peak) amplitude values of the four acquisition heights in six calibration surfaces. The (dotted) 500-m  
258 data deviated in two cases more than 5%, which could be due to strip-level variations of the low output power.

## 259 2.5 Evaluation of the receiver's compliance with the radar equation

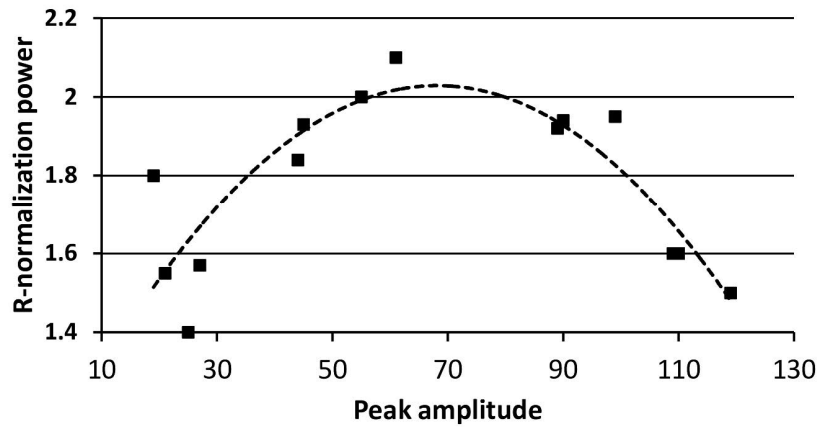
260 Concerning the linearity of the ALS50/60 sensors, Korpela et al. (2010) compared the intensity data of Optech ALTM3100 and  
261 Leica ALS50-ii sensors to assess the influence of vegetation structure in range correction, i.e., the value of  $pwr$  in the range  
262 normalization coefficient  $(R/R_{ref})^{pwr}$ . The results with ALTM3100 were in line with those by Gatzolis (2011) and the optimal  $pwr$   
263 term was 2.1–2.8 for trees and could be logically linked to species-specific structure and echo type. Thus, the  $pwr$  term was  
264 close to 2 for 'only' echoes that are strong and was highest, nearly 2.8, for 'first-of-many' echoes in Scots pine that have the  
265 most 'diffuse' crowns. However, the ALS50-ii intensity data did not comply with the radar equation, and the optima were

266 unclear with the pwr term even below 2. A simple linear model with zero offset was used, which for ALS50-ii (which is similar to  
267 ALS60) is now known to be incorrect (Fig. 6), while in ALTM3100 the zero intensity was closer to the no-signal level.  
268 The amplitude and intensity values in ALS60 were stored as 8-bit signed integers. 'No-signal sections' of the WFs, i.e., the front  
269 and tail parts of the 256-sample-long WFs, showed variance, which also responded to the AGC. The no-signal amplitude values  
270 were about 10–12, while the first-return intensity data was 'truncated' at zero (Fig. 6). The zero point (offset) of the first-return  
271 intensity was negative, and was estimated to be about -24 using Eq. (2), while the offset of A was 11.4 (Section 2.4). The  
272 estimation of the offset of the intensity scale was almost ill-posed, probably because the linear model (response to at-sensor  
273 power) is inaccurate, which is also indicated by Fig. 6.

274 The 2013\_## datasets were acquired to verify the linear response of the ALS60 receiver. In these data, the receiver gain and  
275 transmitted power were fixed. The pwr term of the R-normalization coefficient was again estimated using surfaces of varying  
276 reflectance. Fig. 7 implies that the peak amplitude data (restricted to pulses that had triggered a DR echo) are in a non-linear  
277 dependence with the power entering the aperture. Fig. 8 shows a power → A mapping that creates the pattern in Fig. 7. It  
278 deviates only slightly from a linear model. The component(s) that cause this modest non-linearity in the fixed receiver gain mode  
279 remain unidentified. Possibilities include the Avalanche photodiode and its transimpedance amplifier, other amplifiers and the  
280 digitizer. Furthermore, the non-linearity of low and high signals (Figs. 6, 7, 8) may be due to different components.

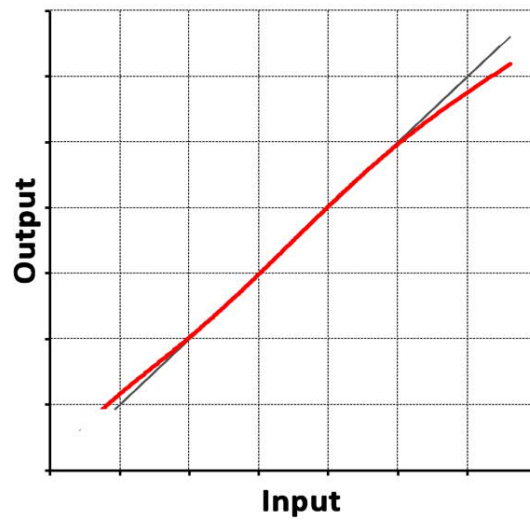


281  
282 Fig. 6. The relationship between first-echo WF peak amplitude and DR intensity in a scene comprising bitumen, tile and metal  
283 roofs, short grass, concrete and asphalt. Echo triggering requires a pA of approximately 20.



284

285 Fig. 7. Estimated power term of range normalization (y-axis) for the pA data (x-axis) in the 2013\_ datasets. In a linear sensor the  
 286 power factor is two. Analyses apply to horizontally aligned surfaces of varying reflectance in pulses that gave rise for an  
 287 observation in 700/800, 700/900 and 800/900 m pulses. The optimal power term minimized the squared sum of normalized  
 288 differences for a particular surface and data pair. The observation is placed at the pA value corresponding to the 800 m  
 289 acquisition.



290

291 Fig. 8. A non-linear sensor response curve that produces the pattern in Fig. 7. The x-axis depicts the power entering the  
 292 aperture, while the y-axis shows the corresponding peak amplitude value.

293 [2.6 Variation of the received waveforms due to sensor effects](#)

294

295 The mean FWHMs in well-defined targets were 10.1–11.3 ns and 6.8–7.9 ns in the 22–59 cm and 11 cm footprint data,  
 296 respectively. The dependence of FWHM on pA was different for the 4 and 9 ns pulses. FWHM increased moderately with pA in  
 297 the 9 ns data, while it decreased considerably with increasing pA in the 4 ns data (Fig. 9).

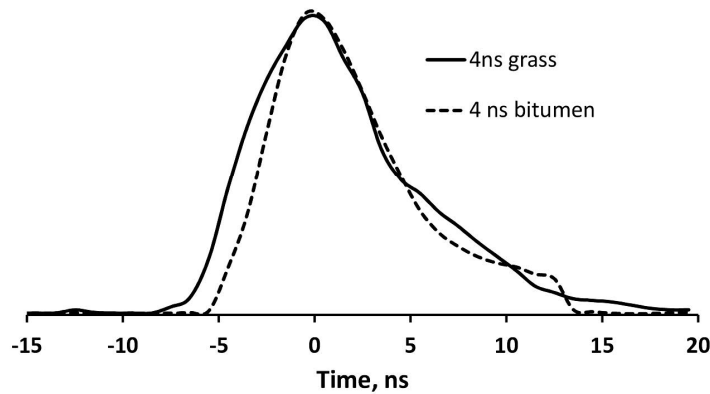
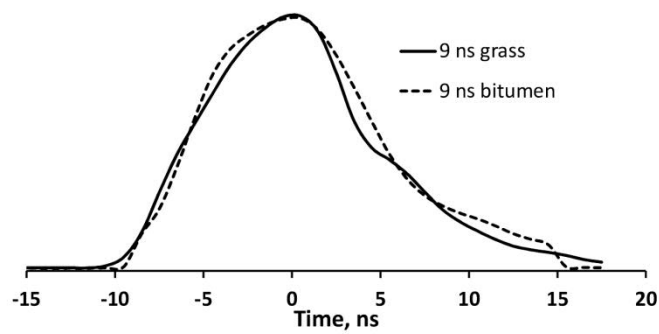


Fig. 9. Averaged, centered and normalized WFs of 4 (7) and 9 (10) ns pulses (transmitted, received). The response of FWHM to surface brightness differs with pulse length. In 4 ns pulses the recorded WF rises in 5–7 ns. The shape difference of the rising part of the 4 ns pulses is caused by the receiver's impulse response, which differs for weak and strong signals. The shape of the 9 ns pulses is retained better in bright and dark targets.

These findings imply that the receiver's response to the fast rising 4 ns pulses differed slightly between weak and strong signals. The CV (coefficient of variation) of FWHM was 1–1.5%. To compare, in some low (300 m) altitude datasets that were acquired in Hyytiälä with the Riegl LMS-Q680i sensor, the FWHM depends strongly on the signal strength, because of the saturation of the amplitude scale due to the low altitude, while 700 m data with the same sensor showed practically no FWHM  $\times$  pA dependence. The echo width should depend only on the target properties, which was not exactly true in the ALS60 data used here.

## 2.7 Evaluation of the radiometric normalization

Radiometric normalization (Eq. 2) was validated across the landscape, using surfaces that were known from the field and/or identified in large-scale aerial images. Each surface was represented by more than 100 pulses per footprint size. The relative differences in Table 3 are normalized with respect to the 11-cm data. The differences were correlated between the 22, 44 and 59 cm data with  $R^2$  of 0.74, 0.86 and 0.89, which implies that the 11 cm data did have offsets due to the within- or between-strip fluctuation of the low (4% of maximum) transmitted power. Using 9% of the maximum would have been at the eye-safety limit. Some of the 11-cm strips were repeated the following day with a 20% (0.8 dB) higher pA. Table 3 also shows that the relative

317 deviations were negatively correlated with the surface brightness, i.e., the pA values in dark surfaces and in the 22–59 cm data  
318 were modestly higher than in the 11 cm data. This can be caused by the different receiver response as the 11-cm data had  
319 shorter pulses with a fast rise in the WF.

320 As stated, the 2012\_## acquisitions unintentionally did not account for atmospheric losses. The at-sensor threshold  
321 backscattering that triggers an echo depends on the instrument, not the acquisition height as such. The atmospheric losses  
322 however caused this threshold to be higher, in terms of the required target backscatter cross-section, in the high-altitude data.  
323 The effect was estimated to be 1.0 dB or 26% between the 500-m and 2700-m data (Section 2.4). Based on the 900-nm field  
324 HCRF measurements of different bitumen surfaces that were carried out in 2008/2009 (Korpela et al. 2011), the minimal echo-  
325 triggering 'reflectance' was about 0.03 in 11 cm data and 0.038 in 59 cm data. For this study I found an even darker bitumen  
326 roof, in which most of the 500 m pulses had produced an echo, while the probability reduced with acquisition height (\*\* in Table  
327 3).

328 This 0–1 dB 'echo-triggering bias' was thus present in all data. It could be observed in comparable (same LiDAR strips, adjacent  
329 targets) data that covered a dark and gently sloping bitumen roof and an adjoining meadow (\*\* in Table 3). In comparison to the  
330 11 cm data, the pA in the 22–59 cm data were 15% lower in bitumen, while the values for meadow were 1–4% higher. The roof  
331 had a reflectance variation due to litter (nearby trees), mosses and lichens (north vs. south exposure). The 11-cm data had more  
332 no-echo pulses, which means that the 11 cm samples were biased towards brighter roof patches, while the larger footprints  
333 averaged the variation. Overall, the validation in different surfaces showed that apart from very dark surfaces, the mean pA was  
334 not affected much by the footprint size (Table 4). Furthermore, the differences in the dark validation surfaces were in all  
335 likelihood caused by small scale reflectance variation. One of the dark surfaces was a pond with dense canopies of sub-emergent  
336 mosses next to a sedge fen. The same scale-related phenomenon was observed here. The mosses did not always give rise to an  
337 echo and the 11-cm targets probably comprise small, dense and drier moss patches.

338

340 Table 3. Relative differences (%) in peak amplitude in different validation surfaces. Values are normalized to the 11 cm data  
 341 except for the first row. Values in which the peak amplitude was lower than in the 11 cm case are positive. Surfaces marked with  
 342 \* or \*\* are adjacent and qualify for pair-wise comparisons. Surfaces #1, #2 and #3 are presented in Section 2.10.

Surface	pA	11	22	44	59
Field	107	–	0.0	-2.2	1.5
	115	0.0	2.1	1.6	3.8
Sedge fen	98	0.0	-4.2	-6.7	-5.2
Low-sedge bog	108	0.0	0.7	-5.8	-4.5
Sedge fen *	107	0.0	-1.9	-1.6	-2.0
Sub-emergent mosses *	24	0.0	12.4	10.9	17.3
Fine sand	87	0.0	0.5	1.4	2.1
Dark bitumen roof **	17	0.0	12.9	15.2	15.1
Meadow **	115	0.0	-4.2	-3.3	-1.2
Forest roads	66	0.0	-2.0	-2.4	-1.9
	77	0.0	0.9	2.0	2.6
#1	80	0.0	0.5	3.0	2.9
#2	70	0.0	3.0	3.2	3.4
	73	0.0	1.1	-4.4	2.2
#3	42	0.0	3.7	3.2	5.9
	50	0.0	3.2	2.5	4.3
	77	0.0	-3.1	-5.7	-5.2
Asphalt roads	51	0.0	4.8	4.2	6.4
	49	0.0	2.5	2.8	4.6
	42	0.0	0.5	0.4	2.9

343

## 344 2.8 Scale-dependent reflectance variation

345

346 Table 4 shows the CV of pA in surfaces of varying small-scale variation. Asphalt and bitumen (different from that in Table 3)  
 347 showed a stable CV. The between-pulse variation reduced with increasing footprint size in hay and in the mire surfaces owing to  
 348 their decimeter-scale spatial variation. The mire was a mosaic of 20–60-cm wide dry hummocks surrounded by wetter white  
 349 mosses and sedges. The multi-footprint data provided logical results. The lowest CV was 2.8%. Bitumen showed relatively high  
 350 CVs, because of the low relative measurement accuracy of weak signals. A comparison between the AGC on (and normalized)  
 351 and AGC off acquisitions could be made with the 2013\_07 dataset. The CVs were slightly lower in the AGC off data: 8.9→8.6%,  
 352 6.9→6.7%, 5.4→4.9% and 4.9→4.0% in bitumen, asphalt, sand and grass, respectively. These results suggest that the AGC



353 normalization was effective and, conditioned on the stability of the targets, that the precision of the instrument was comparable  
354 between 2012 and 2013.

355 Table 4. Coefficient of variation (CV, %) of peak amplitude in different man-made and natural surfaces. Surfaces are in dec-  
356 reasing order of brightness (About 0.45→0.04 in reflectance).

Surface	Footprint, cm			
	11	22	44	59
Hay, < 50 cm	7.0	5.8	4.6	3.6
Short grass	4.9	4.3	4.0	3.5
Mire surface	9.7	8.3	6.3	5.4
Fine sand	4.5	3.5	3.0	2.8
Asphalt, old	7.4	7.3	6.9	7.2
Asphalt, new	5.1	4.9	5.0	5.1
Gravel road	11.4	9.7	7.7	6.9
Bitumen	9.4	9.4	9.0	8.6

357

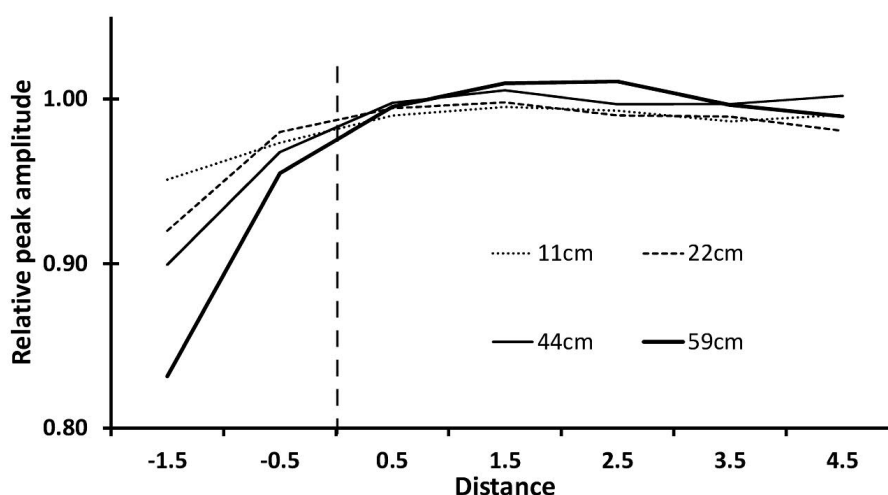
## 358 2.9 Transmission losses in wires and tree crowns

359

360 Next, the data was evaluated using careful measurements of transmission losses caused by three parallel (1.2 m separation, 152  
361 m in total) power line cables 11 m above a spring barley field. The (bright) insulators in the poles were positioned using forward  
362 ray-intersection point estimates obtained in large-scale aerial images and the LiDAR points and constrained by the 1.2 m  
363 separation. Line-to-line distances were used to find pulses that had potentially intersected a cable. The limiting distance was the  
364 95% point of the footprint diameter, i.e. 0.83, 0.62, 0.31 and 0.16 m for the 59, 44, 22 and 11 cm footprint diameters,  
365 respectively. Assuming that the pulse (intersection, 3D LiDAR positions were available) geometry was correct, that the irradiance  
366 profile of pulse cross-section is Gaussian, that the intersection probabilities were equal and that the cables had a 1.2 m  
367 separation, the theoretical losses for a cable of 1 cm thickness are 2.8, 1.5, 0.7 and 0.8%. The observed relative differences of pA  
368 were 2.6, 1.6, 0.7 and 1.1% for the footprint sizes of 11, 22, 44 and 59 cm. The results were thus in line with expectations.

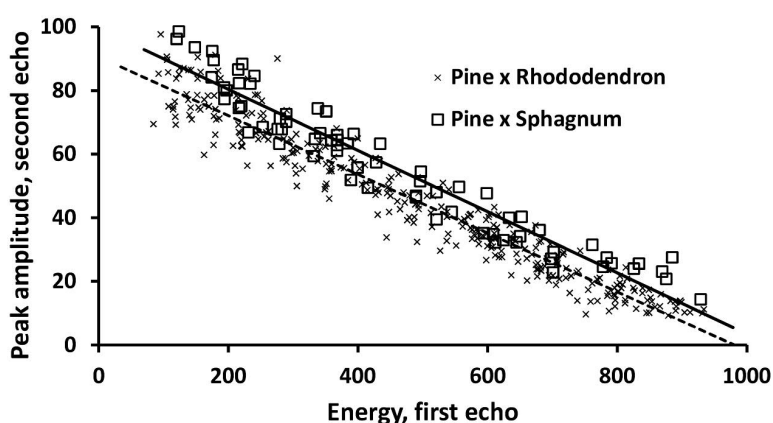
369 The transmission losses were next examined in an approximately 100-year-old barren pine stand having one dominant canopy  
370 layer only. Fig. 10 shows results of relative transmission losses measured by the peak amplitude of ground returns (of single  
371 return pulses) as a function of the 2D trunk-pulse distance at the 60% relative tree height, which approximates crown base  
372 height rather well in this stand. The crown width estimates (N=213) were obtained by fitting a surface of revolution to the LiDAR  
373 point cloud as in (Korpela et al., 2011; Hovi et al., 2015). The ground flora comprises three moss species and patches of reindeer  
374 lichens. Tree heights are 15.9–26.5 m. In Korpela (2008), the ground flora in this stand was mapped using a 960 m<sup>2</sup> photo mosaic  
375 and no correlation was found in the spatial distribution of the bottom flora and the trees, i.e., the reflectance properties of the

376 forest floor are spatially stationary. The pA of ground returns of pulses that had passed through the crown decreased with  
 377 increasing footprint size (Fig. 10). Thus, less energy reached the ground when the footprint became larger and the single ground  
 378 echo was weaker. Fig. 11 shows how the losses, as measured by energy of the first echo, correlate negatively with the peak  
 379 amplitude of the ground return in two pine bogs. The other site has Sphagnum species and tussock cottongrass in the  
 380 understory, while the ground flora of the other bog comprises a dense canopy of 30–60-cm-high wild rosemary. The reflectance  
 381 differences of the low vegetation are reflected by the peak-amplitude values of the second echoes.



382

383 Fig. 10. Relative peak amplitude of (single) ground returns below a Scots pine canopy as a function of the pulse-trunk distance  
 384 (m) at the crown base height. The amplitudes were normalized to data 3–5 m away from the trunks. A distance of zero  
 385 corresponds to the case in which the pulse intersected the crown perimeter at the crown base height. A negative value means  
 386 that the intersection was inside the crown model.



387

388 Fig. 11. Joint distributions of the energy of the first echo and the peak amplitude of the second echo in 9 ns pulses (1, 2 and 2.7  
 389 km data) from two barren pine bogs that differ in understory flora. The pulses had a single WF peak in each of the two echoes.

390

391

## 392 2.10 WF features in single-species tree canopies

393

394 Next, the response of WF attributes to the varying footprint diameter was evaluated in tree canopies. Three 30-year-old planted,  
 395 non-thinned, pure and dense stands of Scots pine, Norway spruce and Silver birch adjacent to a forest road were chosen. The  
 396 road surface was essential for controlling between-dataset offsets in pA. The roads are marked as #1, #2 and #3 in Table 3. The  
 397 small relative offsets of pA were eliminated from the pA and E differences in Table 5. The energy (E) attribute varied least with  
 398 footprint size. The relative variance was also lowest in E with CVs ranging from 10 to 20%. Interestingly, E increased by 8, 6 and  
 399 7% with footprint size in pine, spruce and birch, respectively. The differences between species were largest in pA and E, which is  
 400 in line with the findings in Hovi et al. (2015). The relative decrease of pA with increasing footprint size is interesting, as it was  
 401 lowest for birch (16.6%), followed by pine (19.4%) and spruce (25.9%). An extra stand of Black alder showed a decrease of  
 402 12.6%, but this stand did not have a control surface. These findings indicate weak species-specific 'signatures'.

403 Table 5. Mean values and CV (%) of first-echo WF attributes in 30-year-old pure pine, spruce and birch stands. The LiDAR plots  
 404 (1300 m<sup>2</sup>) contained approximately 70 trees. FWHMs were normalized to the mean values from well-defined surfaces. The E  
 405 attributes of 11 cm (4 ns pulses) were multiplied by 1.307, which is the average ratio between 11 cm (4 ns) and 9 ns data in well-  
 406 defined surfaces of varying reflectance.

Species	Footprint	FWHM	RiS	pA	Energy
Birch	11	1.51 (38)	1.35 (28)	66.1 (34)	1141 (19)
	22	1.49 (38)	1.30 (25)	67.9 (30)	1163 (14)
	44	1.68 (40)	1.51 (26)	58.7 (29)	1205 (13)
	59	1.77 (39)	1.63 (25)	55.2 (28)	1235 (12)
Spruce	11	1.33 (40)	1.21 (26)	60.2 (33)	941 (20)
	22	1.40 (41)	1.25 (28)	57.2 (32)	935 (16)
	44	1.66 (47)	1.45 (30)	49.6 (34)	982 (15)
	59	1.84 (47)	1.60 (30)	44.6 (33)	995 (15)
Pine	11	1.45 (37)	1.30 (27)	53.4 (30)	869 (16)
	22	1.49 (35)	1.32 (24)	51.8 (25)	875 (12)
	44	1.68 (34)	1.53 (23)	46.1 (22)	916 (10)
	59	1.77 (37)	1.63 (23)	43.1 (25)	932 (13)

407

408 FWHM increased as footprint diameter increased, which is according to expectations. When the footprint size increased from 11  
 409 to 59 cm the relative increase in FWHM was 17% in birch, 38% in spruce and 22% in pine. The rise time attribute showed a  
 410 similar pattern with smaller differences between species, compared to echo width, i.e. FWHM. Depending on the range  
 411 detection algorithm, the increase of rise time can influence the height distribution of canopy echoes such that the wide footprint

412 points have lower heights. This phenomenon was not tested here, because it would have required the implementation of a  
413 range detection algorithm and careful geometric calibration of the LiDAR data. It can also be seen that while  $\rho_A$  decreases  
414 (dampens) with footprint size,  $E$  increases moderately as return pulses stretch.

### 415 3. Discussion

416  
417 This study was a largely successful attempt to acquire multi-footprint LiDAR data in which the SNR was fixed for well-defined  
418 surfaces, and to carry out radiometric analyses with these data. Basic hypotheses concerning the influence of footprint diameter  
419 were verified, although comprehensive testing in tree canopies was left to the future. The datasets that have been acquired are  
420 unique as are the investigations performed using those data. [Hopkinson \(2007\)](#) studied the influence of footprint size and  
421 acquired dual-footprint data using a sensor in which the divergence could be altered. He reports that while all other parameters  
422 remained constant, the intensity values, instead of not changing in well-defined surfaces or changing by a factor of 0.38 in linear  
423 targets, reduced by a factor of 0.22 when the divergence was changed from 0.3 to 0.8 mrad.

424 The ALS60 sensor had a fixed beam divergence, but owing to the adjustable output power, footprint diameters ranging from 11  
425 to 59 cm could be realized. The used power level was a compromise as over 50% more could have been used. The use of full  
426 power would have increased the signals by 3 dB, but in that scenario, only 11 dB of power adjustment margin would have been  
427 available and the maximum acquisition height would have been less than 2 km. Furthermore, it was recommended to reduce  
428 the power of the lowest acquisition height.

429 Unexpectedly the campaign planning software did not account for atmospheric losses and a 1 dB trend remained in the data. It  
430 caused the minimal echo-triggering target backscatter cross-section to increase with acquisition height and footprint size. A  
431 fairly realistic estimate of the atmospheric losses, 0.22 dB/km, was derived as a sub product of the radiometric normalization.  
432 Yet, it could not be ruled out that the observed trend was partially caused by systematic errors in the adjustment of the  
433 transmitted power.

434 Backscatter reflectance or peak amplitude normalization was done using natural surfaces assuming that the peak amplitude  
435 observations are in a linear relationship with the peak at-sensor power (irradiance). The linearity assumption had to be  
436 validated, using a separate multi-height LiDAR dataset, which was acquired for that purpose. The validation was based on the  
437 dependence between the power received and the spherical losses from well-defined surfaces. Fixed pulse power and receiver  
438 gain were applied and the signals from well-defined surfaces were, according to the radar equation, 1.6 (2 dB) times stronger  
439 from 700 m compared to 900 m. Setting the output power and receiver gain correctly was crucial. While operating, the AGC  
440 circuit in the receiver can adjust and compensate the internal signal levels by up to 3 dB. The brightest targets had pulses that  
441 were saturated and distorted in the 700 m DR intensity data. These targets were omitted from the calibration ([Fig. 7](#)), because it

442 remained unclear if the saturation applied to the WFs as well. For example, dense canopies of (planophile) raspberry showed  
443 saturated intensity data. Because of the 2-dB difference in power received, the darkest objects that still triggered an echo in the  
444 700 m data would not show in the 900 m data, and the brightest targets were rejected as the dynamic range of the receiver was  
445 exceeded in the 700 m acquisition. Natural and man-made surfaces were used in testing the linear response. The theoretical  
446 differences in power received are 23.4% and 26.5% between 700 m vs. 800 m and 800 vs. 900 m, respectively. The within-  
447 surface CVs of peak amplitude ranged from 3 to 15%, which means that the estimation of the power term of the range  
448 normalization coefficient was susceptible to random errors. A 2% error the mean peak amplitude (of a surface, e.g. in 800 m  
449 data) results in a coefficient estimate that is off by 0.2. The influence of the incidence angle, or rather, the scan zenith angle, was  
450 weak and could be omitted in all analyses.

451 The resulting normalized peak amplitude data (Eq. 2) could be described as interval-scale observations of backscatter cross-  
452 section or reflectance in well-defined targets, at 1064 nm. The smallest values corresponded to a reflectance of about 0.03,  
453 which was observed in one very dark bitumen roof. This estimate was based on previous field measurements of HCRF carried  
454 out for research in the reflectance calibration of aerial images (Korpela et al., 2011). If full power had been applied (eye-safety, 3  
455 dB higher), the echo-triggering reflectance would have been approximately 0.015–0.02. In Korpela et al. (2013) it was found,  
456 using the 2012\_10 (22 cm) data, that the minimum relative silhouette area in pine and spruce branches that produced an echo  
457 was about 10%. Fig. 2 suggests that the ratio of mean peak amplitude in bitumen and pine crowns is about 2.5 and the 900 nm  
458 HCRF reference for the bitumen roof in Fig. 2 was 0.05. If the reflectance of pine needles is 0.3–0.4, the relative silhouette of 0.1  
459 corresponds to a ‘well-defined-surface reflectance’ of about 0.03–0.04. These calculations are coarse and inaccurate, but  
460 suggest that some coherence exists between earlier findings in canopies and those that were made here about well-defined  
461 surfaces.

462 The observations were shown to be influenced by instability of the transmitted power, in the low-power, low-altitude, small-  
463 footprint data. The sensor manufacturer also pointed out this possibility and the repeated 500 m acquisitions in the following  
464 day also suggested this. It remained unclear what kind of temporal patterns these variations had, but at least a campaign-level  
465 offset was found. It may well be that the power fluctuations are smaller in sensors that apply a few fixed power levels and dual-  
466 channel receivers.

467 The impulse response and bandwidth characteristics were analyzed and only minor signal-strength dependent deformations of  
468 the WFs were observed in the 4 ns pulses. The deformations were neglected in the analyses here. Our experience with the Riegl  
469 LMS-Q680i sensor (cf. Armston et al., 2010) is worth noting here. The output power of LMS-Q680i has fixed levels and is tightly  
470 coupled with the pulse repetition frequency (PRF). Some low altitude (300 m) WF data were ‘overexposed’ and mapped to the  
471 non-linear part of the amplitude scale, which created an unwanted dependence of the echo width on signal strength, which was  
472 not present in a 700-m dataset acquired using the same PRF. LMS-Q680i has two receivers, and WF samples are stored from

473 both receivers when necessary. Such dual-receiver, high- and low-gain designs are common and for example the new Leica  
474 ALS80-series sensors employ such a design. The gain difference of the channels is much larger than the 3 dB gain range of the  
475 AGC in ALS60. Compared with having two gain levels, the AGC in ALS60 could be considered a nuisance. In all, managing the  
476 wide dynamic range of LiDAR data acquisitions calls for different solutions, which are implemented in the transmitter and/or in  
477 the receiver. It would be better for radiometric analyses of LiDAR data if the sensor properties, i.e., an absolute calibration, were  
478 known. The reverse-engineering type of analyses that were carried out in this study would then not be needed for carrying out  
479 radiometrically quantitative LiDAR remote sensing.

480 The presented analyses revealed that the peak amplitude data in ALS60 was not in a very exact linear relationship with the  
481 power entering the aperture. However, this finding was obtained with the AGC-amplifier set at fixed gain and it cannot be ruled  
482 out that the AGC amplifier, when operating, causes further problems. However, this is unlikely, because the BW of such  
483 amplifiers can be expected to be larger than the BW of the digitizers or the photon detector. In [Hovi et al. \(2015\)](#), no evidence  
484 was found of systematic differences in peak amplitude or WF shape that were due to the two oscilloscopes behaving differently.  
485 As said, the oscilloscopes take turns in processing pulse pairs into WF data.

486 The dependence of 'unregistered' transmission losses on footprint size, preceding the triggering of waveform recording, was  
487 shown in power line cables and in pine crowns. The losses were from 2.6 to 0.5% in the cables and from 5 to 16% in the trees,  
488 for the footprint diameters of 11–59 cm. Canopies are comprised of gaps and clumped scatterers, while the cable is a single  
489 reflector obstructing the energy from reaching the ground, where the losses were detected as reductions in the average peak  
490 amplitude of hundreds of pulses. The sites used for loss detection had low and homogenous or at least spatially stationary  
491 ground reflectance variation. The influence of canopy transmission losses can be compensated to some degree in discrete-  
492 return ([Korpela et al., 2012](#)) and WF data (e.g. [Richter et al., 2015](#)) by imposing constraints such as constant reflectance or  
493 configuration of needles in a tree crown, which reduces the ambiguities related to backscatter cross-section. This was  
494 demonstrated by the comparison of the two bogs

495 The relative variation of peak amplitude was assessed in surfaces of varying small-scale structure ([Section 2.8](#)). Fine sand, which  
496 has a HCRF of about 0.2 at 1064 nm, showed the lowest CV, 4.5% in 11-cm data, and as low as 2.8% in the 59-cm data. The  
497 target was a volleyball field with a 'moderately bumpy' surface and a few dry leaves blended in the sand. Bitumen was least  
498 influenced by footprint size, but the CVs were higher, on the order of 9%, because of the low reflectance. The results were  
499 logical, concerning the influence of the footprint size and the averaging effect in large footprint data. The observed CVs suggest  
500 that the instrument is rather stable or precise as there were always data from two or more strips per test surface and footprint  
501 size.

502 Regarding the influence of the footprint size on the tested first-echo WF attributes in tree canopies, there were minor species-  
503 specific differences or 'signatures'. For example, the peak amplitude decreased most in Scots pine stand, the structure of which

504 is less clumped or 'more diffuse' compared to Norway spruce. The reduction in peak amplitude with increasing footprint size  
505 was smallest in the case of birch, which have compact crowns at the age of 30 years. Actually these findings that relate to  
506 canopy structure are line with those concerning the response of intensity data to R variation (Korpela et al., 2010a; Gatsiolis,  
507 2011), because the decrease of power received that is due to an increase of divergence is also dependent on the structure  
508 (Section 1). As hypothesized, the echo width increased with footprint size and the relative increase was largest in the canopy of  
509 spruce, which is reasonable, given the 'clumped structure' and 'layered compact branches'. The energy was the largest in birch,  
510 and this could possibly be due to the slightly higher reflectance of the leaves. Then again, backscattering in birch is strongly  
511 influenced by leaf arrangement as well (Hovi and Korpela, 2014). The tests in forest canopies were hampered by the remaining  
512 between-dataset SNR differences, which were subtracted using the roads adjacent to the test stands as control surfaces.

#### 513 4. Conclusions

514  
515 This study showed that it is possible to acquire multi-footprint data that is radiometrically comparable for well-defined surfaces,  
516 using a commercial topographic LiDAR sensor. The study has shown how to carry out careful analyses to verify that the WF  
517 amplitude observations are on the interval scale. The impulse response and bandwidth characteristics of the sensor were tested  
518 in a simple manner to verify that the WF features depend only on the target-pulse interaction, in order to draw correct  
519 conclusions from observations concerning the influence of the footprint size. Many of the outlined hypotheses regarding the  
520 influence of footprint size could be confirmed in real data. These included the widening of the echoes, the increase of return  
521 energy, and the decrease of maximal amplitude in tree canopies, all with increasing footprint size. These effects had species-  
522 specific traits. Transmission losses as small as 0.5% could be detected by an appropriate selection of the test site and by using a  
523 large amount of pulses to cancel out the random variation. In conclusion, multi-footprint data may offer some advances in the  
524 estimation of tree or forest canopy structure. However, the species-specific differences are likely quite weak. A thorough  
525 investigation is needed to see if multi-footprint LiDAR can enhance forest remote sensing. For example, the estimation of  
526 canopy closure might benefit from dual-footprint data. The experience that was gained here constitutes a good starting point.

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535 Armston, J., Disney, M., Lewis, P., Scarth, P., Phinn, S., Lucas, R., Bunting, P. & Goodwin, N. (2013). Direct retrieval of canopy gap  
536 probability using airborne waveform lidar. *Remote Sensing of Environment* 134:24–38.

537 Disney, M.I., Kalogirou, V., Lewis, P., Prieto-Blanco, A., Hancock, S. & Pfeifer, M. (2010) Simulating the impact of discrete-return  
538 lidar system and survey characteristics over young conifer and broadleaf forests, *Remote Sensing of Environment* 114(7):  
539 1546–1560.

540 Gatzliolis, D. (2011). Dynamic range-based intensity normalization for airborne, discrete return LiDAR data for forest canopies.  
541 *Photogrammetric Engineering & Remote Sensing*. 77(3): 251–259.

542 Goodwin, N.R., Coops, N.C. & Culvenor, D.C. (2006). Assessment of forest structure with airborne LiDAR and the effects of  
543 platform altitude. *Remote Sensing of Environment* 103(2): 140–152.

544 Hopkinson, C. (2007). The influence of flying altitude, beam divergence, and pulse repetition frequency on laser pulse return  
545 intensity and canopy frequency distribution. *Canadian Journal of Remote Sensing* 33: 312–324.

546 Hovi, A. & Korpela, I. (2014). Real and simulated waveform-recording LiDAR data in juvenile boreal forest vegetation. *Remote*  
547 *Sensing of Environment* 140: 665–678.

548 Hovi A., Korhonen L., Vauhkonen J. & Korpela I. (2015). LiDAR waveform features for tree species classification and their  
549 sensitivity to tree- and acquisition related parameters. *Remote Sensing of Environment* 173: 224–237.

550 Hovi, A. (2015). Towards an enhanced understanding of airborne LiDAR measurements of forest vegetation. *Dissertationes*  
551 *Forestales* 200. 69 p.

552 Korpela, I. (2008). Mapping of understory lichens with airborne discrete-return LiDAR data. *Remote Sensing of Environment*  
553 112(10): 3891-3897.

554 Korpela, I., Ørka H.O., Heikkinen V., Tokola T., & Hyypä J. (2010a). Range- and AGC normalization of LIDAR intensity data for  
555 vegetation classification. *ISPRS Journal of Photogrammetry and Remote Sensing* 65(4): 369–379.

556 Korpela, I., Ørka, H.O., Maltamo, M., Tokola, T. & Hyypä, J. (2010b). Tree species classification using airborne LiDAR – effects of  
557 stand and tree parameters, downsizing of training set, intensity normalization and sensor type. *Silva Fennica* 44(2): 319–339

558 Korpela, I., Heikkinen, V., Honkavaara, E., Rohrbach F. & Tokola, T. (2011). Variation and anisotropy of reflectance of forest trees  
559 in radiometrically calibrated airborne line sensor images – implications for species classification in digital aerial images. *Remote*  
560 *Sensing of Environment* 115(8): 2062–2074.



561 Korpela I., Hovi A. & Morsdorf F. (2012). Understory trees in airborne LiDAR data - Selective mapping due to transmission losses  
562 and echo-triggering mechanisms. *Remote Sensing of Environment* 119: 92–104.

563 Korpela I., Hovi A. & Korhonen L. (2013). Backscattering of individual LiDAR pulses explained by photogrammetrically derived  
564 vegetation structure. *ISPRS journal of photogrammetry and remote Sensing* 83: 81–93.

565 Maltamo M. & Packalén P. (2014). Species-specific management inventory in Finland. In: Maltamo M., Næsset E., Vauhkonen J.  
566 (eds.). *Forestry applications of airborne laser scanning: Concepts and case studies*. Springer, The Netherlands, p. 241–252

567 Mallet, C. & Bretar, F. (2009). Full-waveform topographic lidar: State-of-the-art, *ISPRS journal of photogrammetry and remote*  
568 *sensing* 64(1):1–16.

569 Ørka, H-O., Næsset, E., Bollandsås, O., M. (2010). Effects of different sensors and leaf-on and leaf-off canopy conditions on echo  
570 distributions and individual tree properties derived from airborne laser scanning. *Remote sensing of environment* 114(7): 1445–  
571 1461.

572 Richter, K., Blaskow, R., Stelling, N., & Maas, H. G. (2015). Reference Value Provision Schemes for Attenuation Correction of Full-  
573 Waveform Airborne Laser Scanner Data. *ISPRS Ann. Photogramm. Remote Sens. Spat. Inf. Sci.* 1: 65–72.

574 Roncat, A., Briese, C., Jansa, J. & Pfeifer, N. (2014). Radiometrically calibrated features of full-waveform lidar point clouds based  
575 on statistical moments. *IEEE Geoscience and remote sensing letters*, 11(2): 549–553

576 Wagner W., Ullrich A., Ducic V., Melzer T., Studnicka N. (2006). Gaussian decomposition and calibration of a novel small-  
577 footprint full-waveform digitising airborne laser scanner. *ISPRS Journal of Photogrammetry and Remote Sensing* 60: 100–112.

578 Wagner, W. (2010). Radiometric calibration of small-footprint full-waveform airborne laser scanner measurements: Basic  
579 physical concepts, *ISPRS Journal of Photogrammetry and Remote Sensing* 65 (6): 505–513.

580 Vauhkonen, J., Maltamo, M., McRoberts, R.E., & Naesset, E. (2014). Introduction to forestry applications of airborne laser  
581 scanning. In: Maltamo M., Næsset E., Vauhkonen J. (eds.). *Forestry applications of airborne laser scanning: Concepts and case*  
582 *studies*. Springer, The Netherlands, p. 1–16.